

Post Processing of Hurricane Model Forecasts

T. N. Krishnamurti

Florida State University

Tallahassee, FL



Collaborators: Anu Simon, Mrinal Biswas, Andrew Martin, Christopher Davis,
Aarolyn Hayes, Naomi Surgi, TSV Vijaya Kumar



ARW Model Description

The real-time ARW forecasts in 2005 used a **two-way nested configuration** ([Michalakes et al. 2005](#)), that featured a 12-km outer fixed domain with a movable nest of **4/1.33-km grid spacing**.

The nest was centered on the location of the minimum 500-hPa geopotential height within a prescribed search radius from the previous position of the vortex center (or within a radius of the first guess, when first starting).

Nest repositioning was calculated every 15 simulation minutes and the width of the search radius was based on the maximum distance the vortex could move at 40 m s^{-1} .

On the 12-km domain, the **Kain–Fritsch cumulus parameterization** was used, but domains with **finer resolution had no parameterization**.

All domains used the WRF single-moment 3-class (WSM3) microphysics scheme ([Hong et al. 2004](#)) that predicted only one cloud variable (water for $T > 0^\circ\text{C}$ and ice for $T < 0^\circ\text{C}$) and one hydrometeor variable, either rainwater or snow (again thresholded on 0°C).

Both domains also used the Yonsei University (YSU) scheme for the planetary boundary layer ([Noh et al. 2003](#)).

This is a first-order closure scheme that is similar in concept to the scheme of [Hong and Pan \(1996\)](#), but appears less biased toward excessive vertical mixing as reported by [Braun and Tao \(2000\)](#).

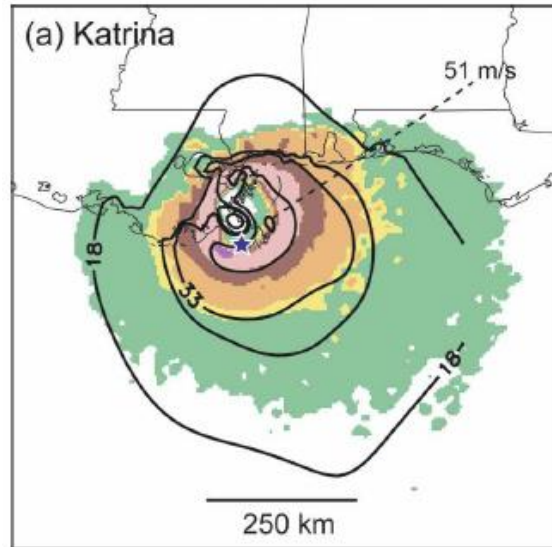
The drag formulation follows [Charnock \(1955\)](#) and is described more in [section 5](#). The surface exchange coefficient for water vapor follows [Carlson and Boland \(1978\)](#), and the heat flux uses a similarity relationship ([Skamarock et al. 2005](#)).

The forecasts were integrated from 0000 UTC and occasionally 1200 UTC during the time when a hurricane threatened landfall within 72 h.

Forecasts were initialized using the Geophysical Fluid Dynamics Laboratory (GFDL) model, with data on a $\frac{1}{2}^\circ$ latitude–longitude grid. The Global Forecast Model (GFS) from the National Centers for Environmental Prediction (NCEP), obtained on a 1° grid, was used only when the GFDL was unavailable.

Post processing diagnostics

Here we shall be showing some post processing for WRF-ARW. This model is presently being added to our suite of mesoscale models.

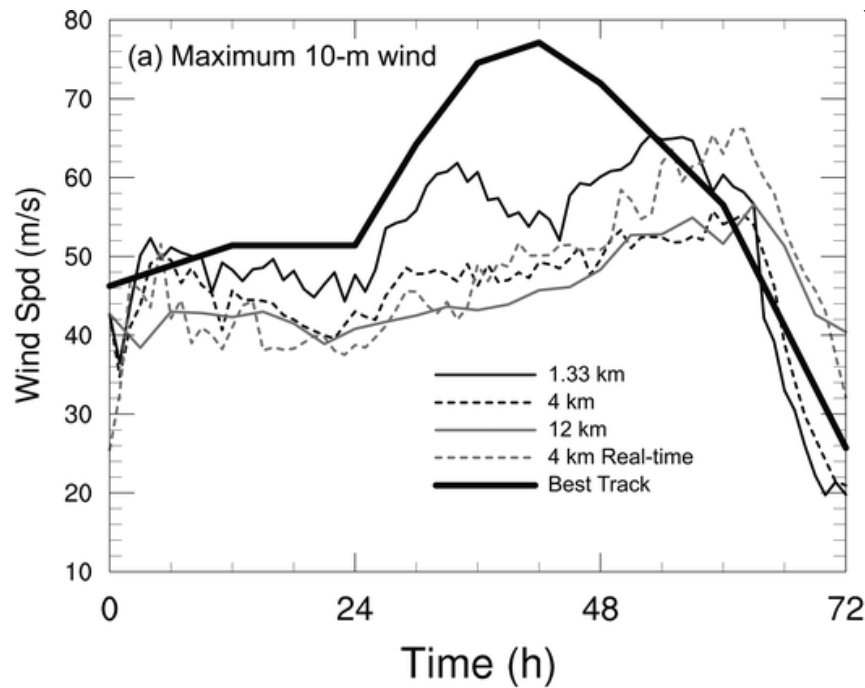


**Katrina, valid time = 1200
UTC 29 Aug (60-h forecast)**

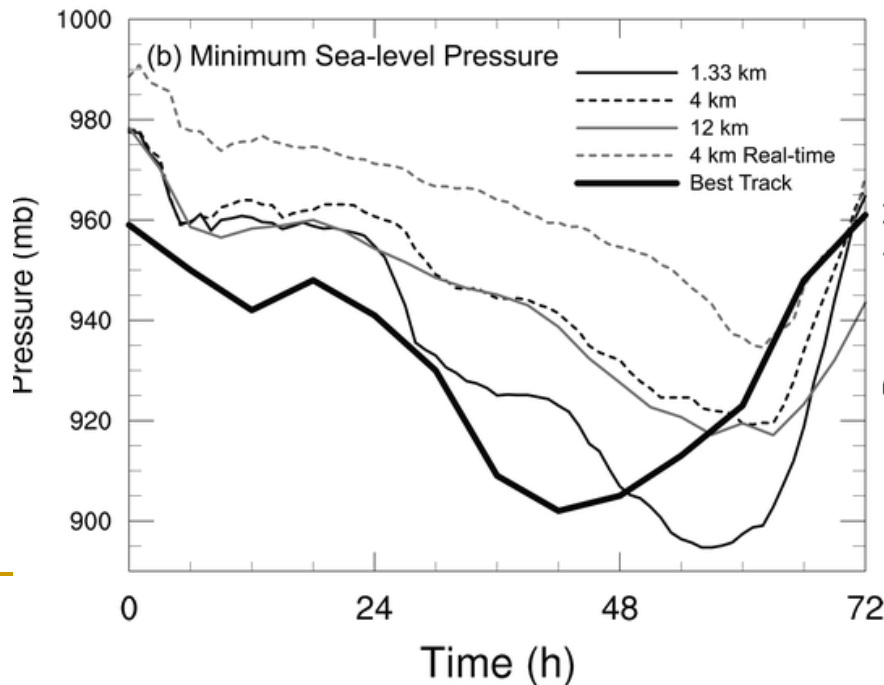
Predicted storm center location at indicated valid times (below) is denoted by blue star in each figure. Wind fields from AHW forecasts have been shifted to observed locations to facilitate comparison.

**HWind valid times are (a) 1132 UTC
29 Aug**

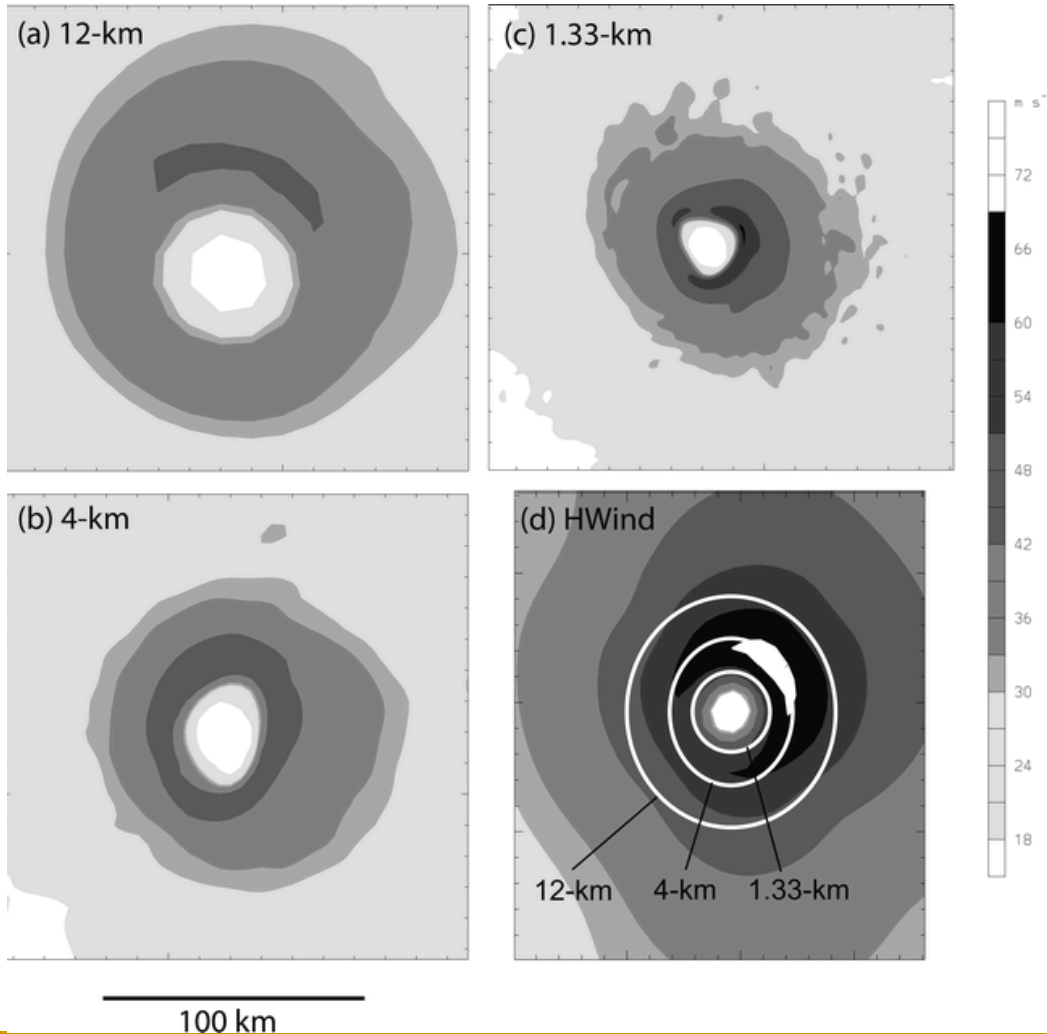
10-m wind from AHW real-time forecasts with contours of nearest HWind (black lines) analyses overlaid



Predicted intensity and minimum sea level pressure at different forecast hours



(a) Maximum 10-m wind and (b) minimum sea level pressure for forecasts of Katrina beginning 0000 UTC 27 Aug. Legend labels 1.33, 4, and 12 km refer to grid spacing of WRF ARW, version 2.1.2, using the Charnock drag relation. The forecast on a 12-km grid used the Kain-Fritsch parameterization. The 4-km real time (gray dashed) refers to the forecast made in real time with an innermost nest of 4-km grid spacing. All retrospective forecasts were initialized with the GFDL initial condition.

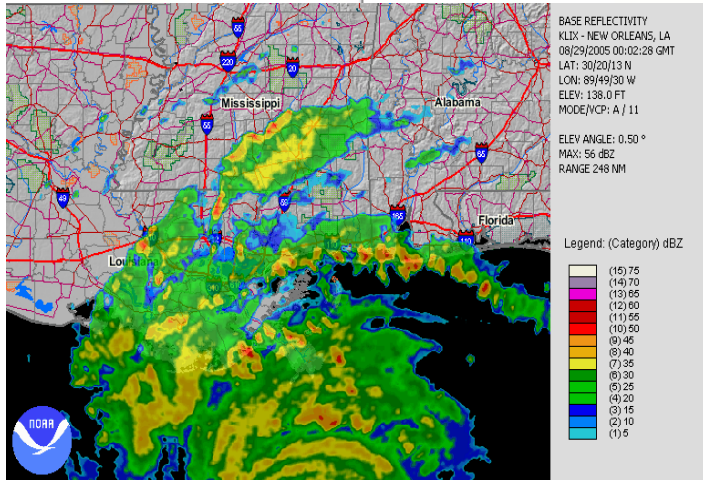
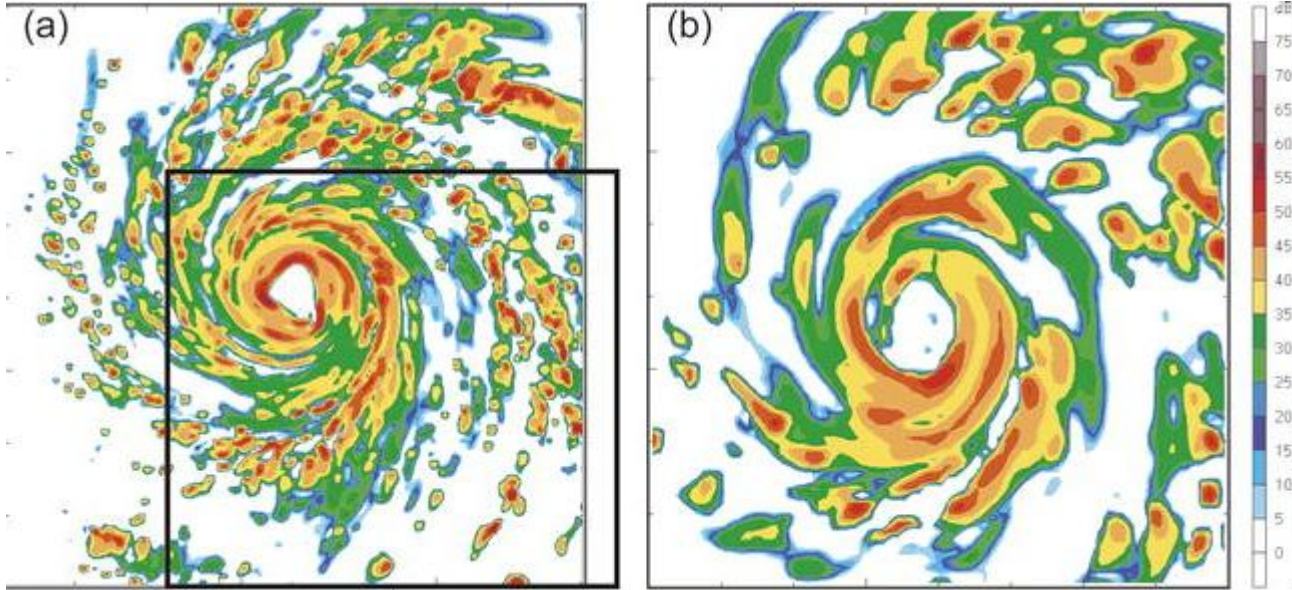


Size of the storm as seen by the predicted wind field at different resolutions compared to HWIND.

Shown here is 10-m wind speed ($m s^{-1}$) from 36-h Katrina forecast valid 1200 UTC 28 Aug on (a) the 12-km grid, (b) the 4-km grid, (c) the 1.33-km grid, and (d) the NOAA HWind product valid 1200 UTC 28 Aug. White ellipses in (d) are an approximate trace of the radii of maximum wind at each azimuth around the vortices in (a), (b), and (c).

1.33 km

4 km



Model-derived reflectivity at 3-km MSL valid 2300 UTC 28 Aug from nest with (a) 1.33-km grid increment and (b) 4-km grid increment. (c) Observed radar reflectivity composite valid between 2000 and 2100 UTC 28 Aug based on tail Doppler radar data from both the NOAA P-3 (red track) and the Naval Research Laboratory P-3 (pink track) with the Electra Doppler radar (ELDORA).

OBS composite radar

Departures from balance laws

The full divergence equation can be written in the form (from Fankhauser 1974):

$$\begin{aligned} -\nabla^2 \phi = & \underbrace{-f \left(\frac{\partial u}{\partial y} - \frac{\partial v}{\partial x} \right) - 2 \left(\frac{\partial v}{\partial x} \frac{\partial u}{\partial y} - \frac{\partial u}{\partial x} \frac{\partial v}{\partial y} \right) + \beta u}_{\text{Red line}} + \underbrace{\left(\frac{\partial \omega}{\partial x} \frac{\partial u}{\partial p} + \frac{\partial \omega}{\partial y} \frac{\partial v}{\partial p} \right) + D^2 +}_{\text{Blue underline}} \\ & \underbrace{\left(\frac{\partial D}{\partial t} + u \frac{\partial D}{\partial x} + v \frac{\partial D}{\partial y} + \omega \frac{\partial D}{\partial p} \right) + \left(\frac{\partial F_u}{\partial x} + \frac{\partial F_v}{\partial y} \right)}_{\text{Blue underline}} \end{aligned}$$

Red lines represent the balance equation (Haltiner and Williams 1980). The blue underlined terms denote the non linear balance which is also expressed as .

$$\nabla^2 \phi = f \nabla^2 \psi + 2J \left(\frac{\partial \psi}{\partial x}, \frac{\partial \psi}{\partial y} \right)$$

Implication of gradient wind departures on hurricane intensity

In local cylindrical storm centered coordinate we can write the complete radial equation of motion in the form:

$$\frac{\partial V_r}{\partial t} + V_\theta \frac{\partial V_r}{r \partial \theta} + V_r \frac{\partial V_r}{\partial r} + \omega \frac{\partial V_r}{\partial p} - \frac{V_\theta^2}{r} - fV_\theta = -g \frac{\partial z}{\partial r} + F_r \text{ or}$$

$$\frac{V_\theta^2}{r} + fV_\theta - g \frac{\partial z}{\partial r} + GWD = 0 \text{ where GWD denotes the gradient wind departure.}$$

$$\text{Where } GWD = -\frac{\partial V_r}{\partial t} - V_\theta \frac{\partial V_r}{r \partial \theta} - V_r \frac{\partial V_r}{\partial r} - \omega \frac{\partial V_r}{\partial p} - F_r$$

$$V_\theta \text{ can be expressed by the relation, } V_\theta = \frac{-f \pm \sqrt{f^2 - \frac{4}{r} \left(GWD - g \frac{\partial z}{\partial r} \right)}}{\frac{2}{r}}$$

This denotes a local value of the tangential wind from a complete radial wind equation in the presence of gradient wind departures GWD. Note that $-g \frac{\partial z}{\partial r}$ is generally <0 in the inner rain area ($r < 200$ km where r is positive outward). The gradient wind, in this locale

$$\text{storm centered coordinate, is given by: } V_\theta = \frac{-f \pm \sqrt{f^2 + \frac{4}{r} \left(g \frac{\partial z}{\partial r} \right)}}{\frac{2}{r}}. \text{ Note that GWD can}$$

be ≥ 0 . Thus several possibilities exist:

If $f^2 < \frac{4}{r} \left(GWD - g \frac{\partial z}{\partial r} \right)$ we have a nonphysical solution.

We are looking for different options for $f^2 - \frac{4}{r} \left(GWD - g \frac{\partial z}{\partial r} \right) > 0$, and its positive root.

Since $g \frac{\partial z}{\partial r} > 0$ we can write the inequality in the form, $f^2 + \frac{4}{r} g \frac{\partial z}{\partial r} > \frac{4}{r} GWD$. The left hand side is essentially positive definite, $GWD < 0$ would always satisfy this condition. A negative value of GWD contributes to a value of V_θ in equation () larger than the gradient wind value, i.e. $V_\theta > V_{g\theta}$, hence the extreme negative values of GWD would go with large values of azimuthal motions i.e. supergradient winds or non physical situations. A positive value of GWD, conversely describes situations with subgradient winds. Such instances of very strong azimuthal motions described by the complete radial equation of motion, are attributed to the departure from gradient wind (GWD) which also relate to large values of divergence, this we illustrate below.

A note on Radial gradient wind solution:

For a hurricane in storm centered coordinate the radial gradient wind equation is expressed by $\frac{V_\theta^2}{r} + fV_\theta = g \frac{\partial z}{\partial r}$ where $\frac{V_\theta^2}{r}$, fV_θ and $g \frac{\partial z}{\partial r}$ are all generally positive. The

roots of this equation are $u = \frac{-f \pm \sqrt{\frac{4}{r} g \frac{\partial z}{\partial r}}}{\frac{2}{r}}$. Note that the negative root of the radical is

nonphysical hence there is only one positive root i.e. $\frac{-f + \sqrt{\frac{4}{r} g \frac{\partial z}{\partial r}}}{\frac{2}{r}}$ for $\left| \sqrt{\frac{4}{r} g \frac{\partial z}{\partial r}} \right| > |f|$.

This implies that $\sqrt{\frac{4}{r} g \frac{\partial z}{\partial r}} > |f|$ applies for cyclonic motions. Note that there are no anomalous solutions of the radial gradient wind equation, i.e. there is only one normal

solution. Such instances of very strong azimuthal motions, described by the complete radical equation of motion, are attributed to the departures from gradient wind (GWD) which relate to large values of divergence.

The full divergence equation can be written in the form (from Fankhauser 1974):

$$-\nabla^2 \phi = -f \left(\frac{\partial V_\theta}{\partial y} - \frac{\partial V_r}{\partial x} \right) - 2 \left(\frac{\partial V_r}{\partial x} \frac{\partial V_\theta}{\partial y} - \frac{\partial V_\theta}{\partial x} \frac{\partial V_r}{\partial y} \right) + \beta u + \left(\frac{\partial \omega}{\partial x} \frac{\partial V_\theta}{\partial p} + \frac{\partial \omega}{\partial y} \frac{\partial V_r}{\partial p} \right) + D^2 +$$

$$\left(\frac{\partial D}{\partial t} + V_\theta \frac{\partial D}{\partial x} + V_r \frac{\partial D}{\partial y} + \omega \frac{\partial D}{\partial p} \right) + \left(\frac{\partial F_{V_\theta}}{\partial x} + \frac{\partial F_{V_r}}{\partial y} \right)$$

Red lines represent the balance equation (Haltiner and Williams 1980). The underlined terms denote the non linear balance which is also expressed as

$$\nabla^2 \phi = f \nabla^2 \psi + 2J \left(\frac{\partial \psi}{\partial x}, \frac{\partial \psi}{\partial y} \right).$$

The complete radial equation of motion in storm centered coordinate (r positive outward) is written in the form:

$$\frac{dV_r}{dt} - \left(f + \frac{V_\theta}{r} \right) V_\theta = - \frac{\partial \phi}{\partial r} + g \frac{\partial \tau_\theta}{\partial p}$$

Or $\frac{\partial V_r}{\partial t} = -V_\theta \frac{\partial V_r}{r \partial \theta} - V_r \frac{\partial V_r}{\partial r} - \omega \frac{\partial V_r}{\partial p} - \frac{V_\theta^2}{r} - f V_\theta = -g \frac{\partial z}{\partial r} + F_r$, the underlined terms

denote radial gradient wind balance. The non linear balance and the radial gradient wind equation are equivalent, Fortak (1956).

Fortak, H., "Concerning the general vertically averaged hydrodynamic equations

Thus departures from non linear balance can be approximately equated to departures from gradient wind balance for the radial direction. Departures from non linear balance largely arise from horizontal and vertical advection of divergence and the divergence square. If a flare up of deep convection occurs near the eye wall of a hurricane divergence (/convergence) increases, so do the departures from balance laws. Growth of negative departures leads to a stronger hurricane. This follows from

$$GWD' = \left(\frac{\partial \omega}{\partial x} \frac{\partial u}{\partial p} + \frac{\partial \omega}{\partial y} \frac{\partial v}{\partial p} \right) + D^2 + \left(\frac{\partial D}{\partial t} + u \frac{\partial D}{\partial x} + v \frac{\partial D}{\partial y} + \omega \frac{\partial D}{\partial p} \right) + \left(\frac{\partial F_u}{\partial x} + \frac{\partial F_v}{\partial y} \right)$$

The complete radial equation may be written as where "GWD" denotes gradient wind departures, and GWD is expressed by the radial wind equations carries the solutions

$$\frac{V_\theta^2}{r} + fV_\theta - g \frac{\partial z}{\partial r} + GWD' = 0$$

the negative root is non physical. The other root, when GWD' is $\ll 0$, carries strong tangential wind for its solution. When $GWD'=0$ then we have a radial gradient wind balance.

$$\frac{V_\theta^2}{r} + fV_\theta = g \frac{\partial z}{\partial r}$$

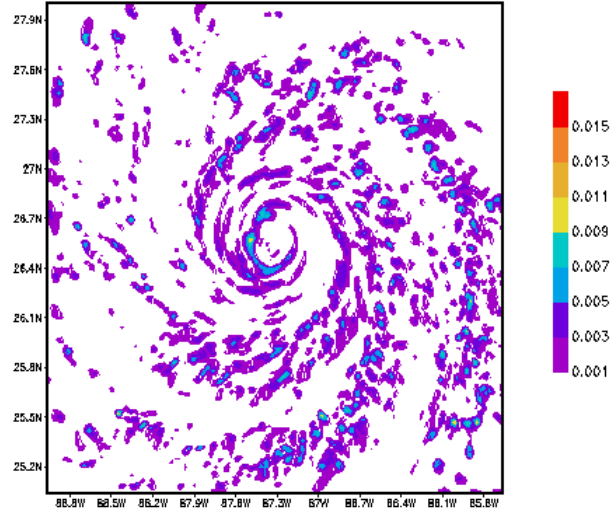
We have routinely mapped the field of GWD' in the intensifying and decaying phases of hurricane intensity.

We have routinely mapped the field of GWD/ in the intensifying and decaying phases of hurricane intensity.

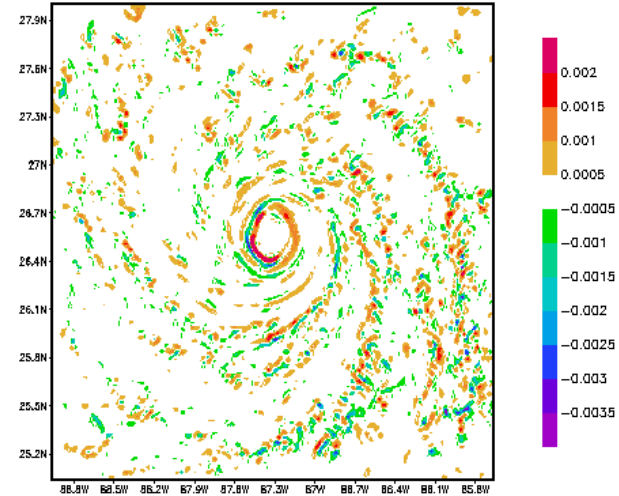
We shall next illustrate several examples of the following scenario:

1. Deep convection flares up near the eye wall, as seen from the local growth of rain water mixing ratio, liquid water mixing ratio or radar reflectivity as implied from model hydrometeors.
 2. Divergence flares up
 3. Departures from balance laws flare up
 4. Solution of complete radial equation shows rapid growth of hurricane intensity.
-

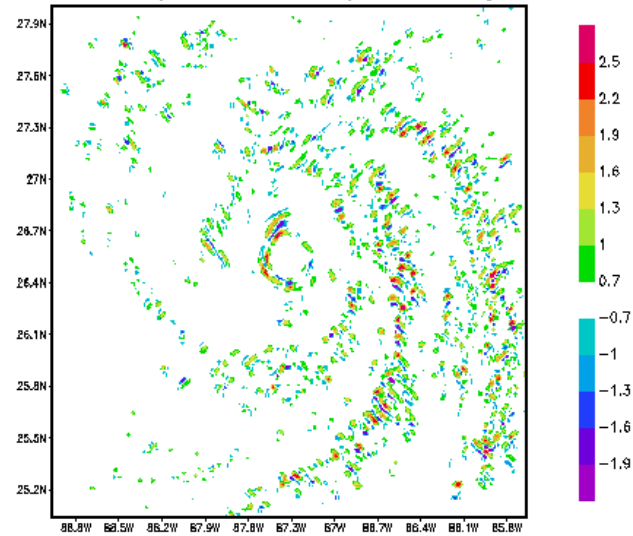
CLOUD LIQUID WATER kg/kg 15z 28 Aug 2005



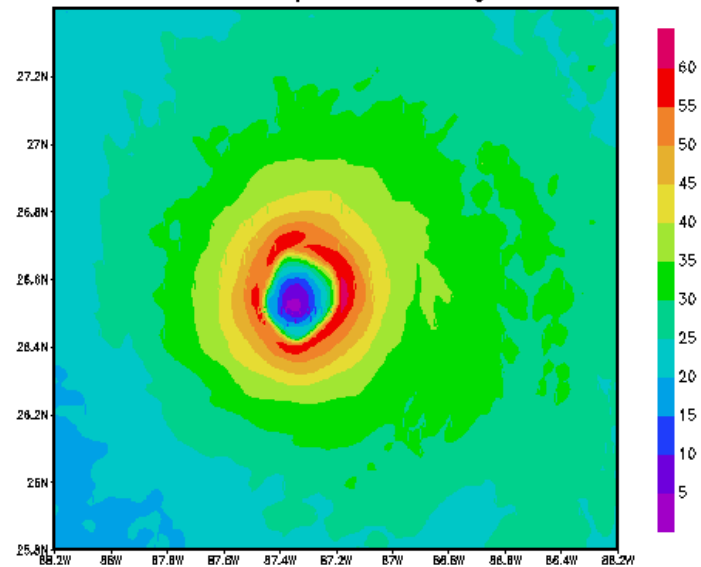
DIVERGENCE (per sec.) 15z 28 Aug 2005



DEPARTURE (x 10⁻⁵sec.**2) 15z 28 Aug 2005

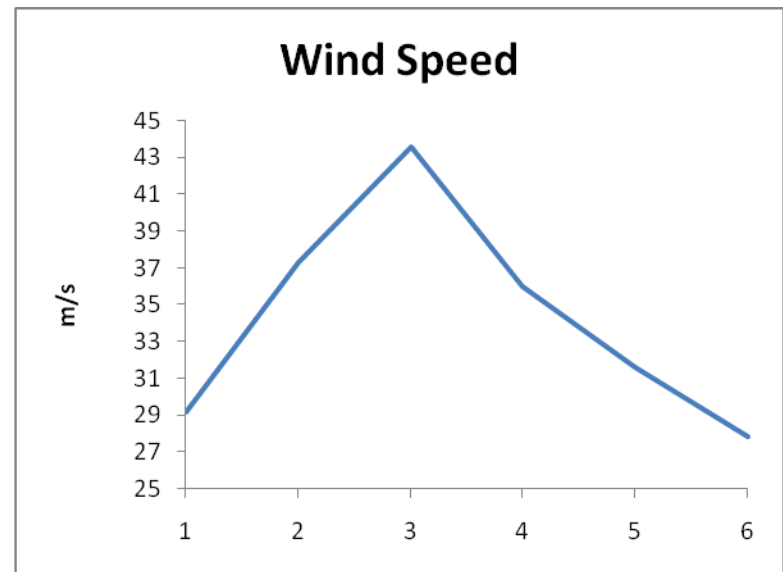
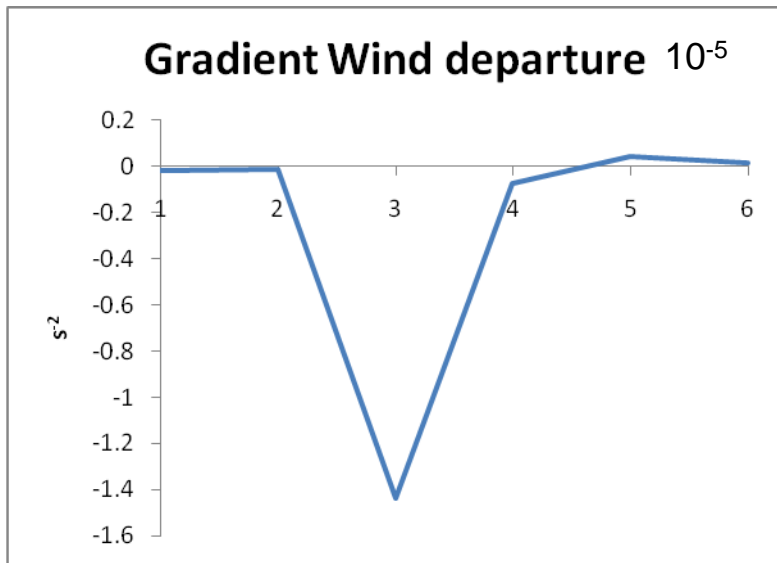
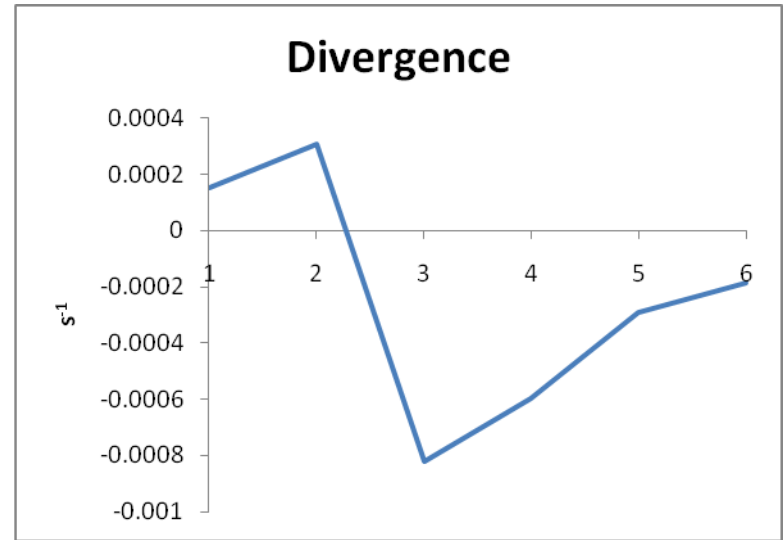
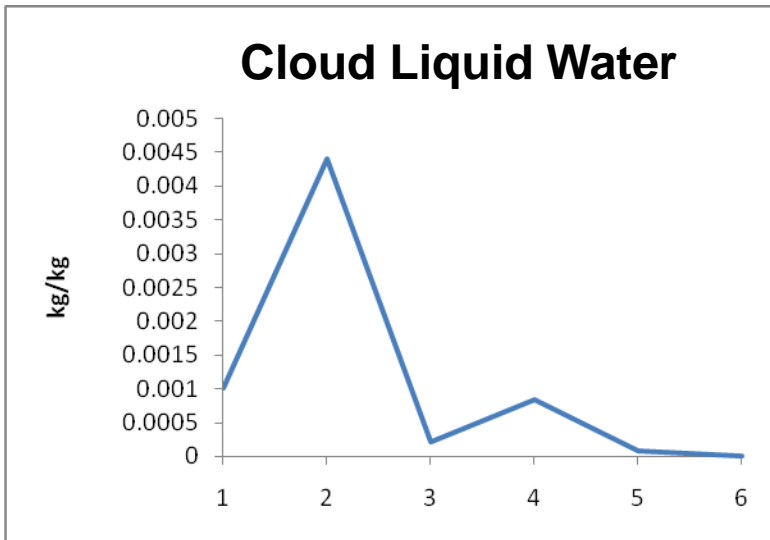


WIND SPEED m/s 15z 28 Aug 2005



These panels correspond closely to the period of maximum intensity of Katrina

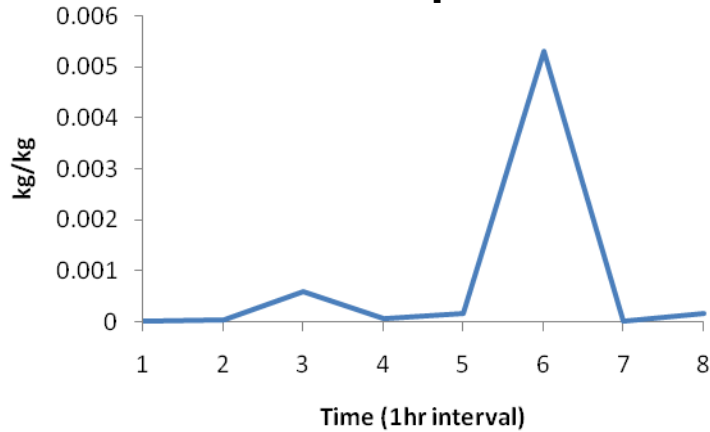
Hourly plots



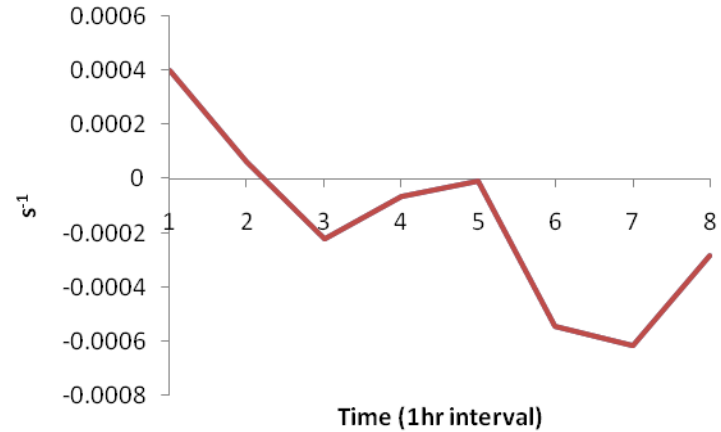
Initial time: 10z 28 August 2005

Hourly plots

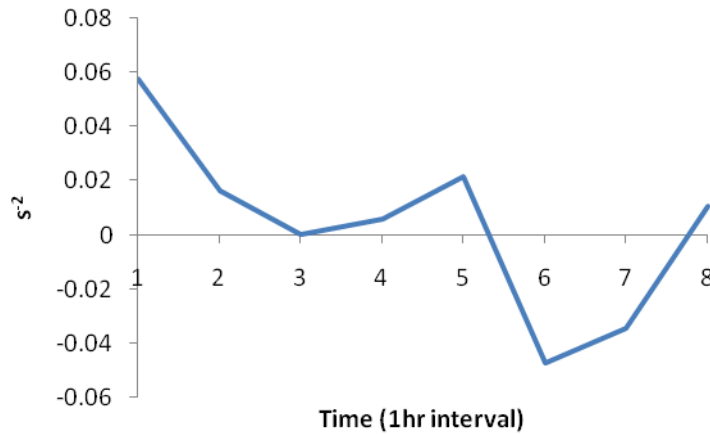
Cloud Liquid Water



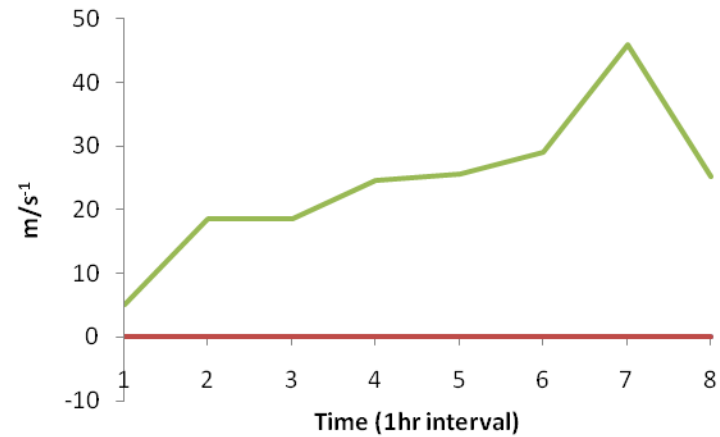
Divergence



Gradient Wind Departure 10^{-5}

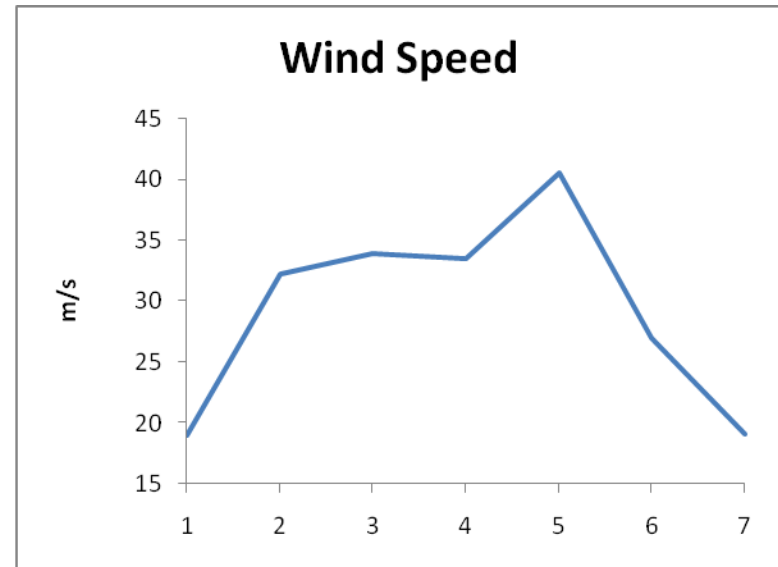
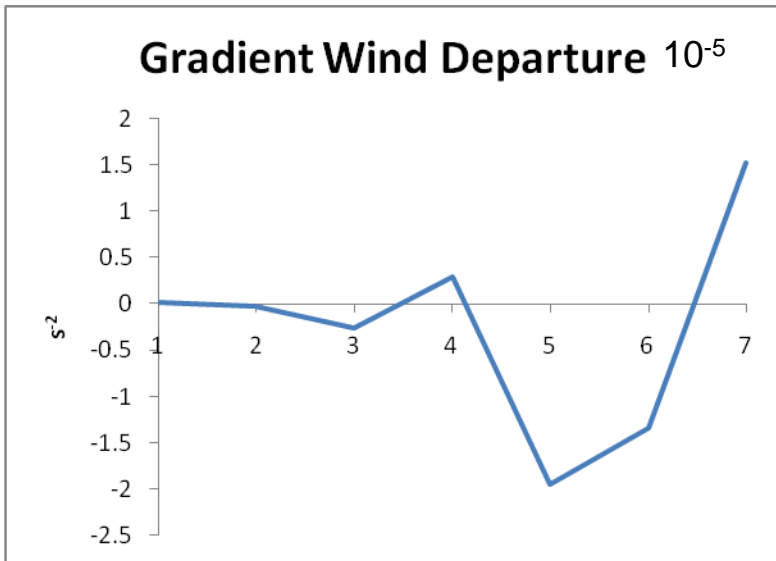
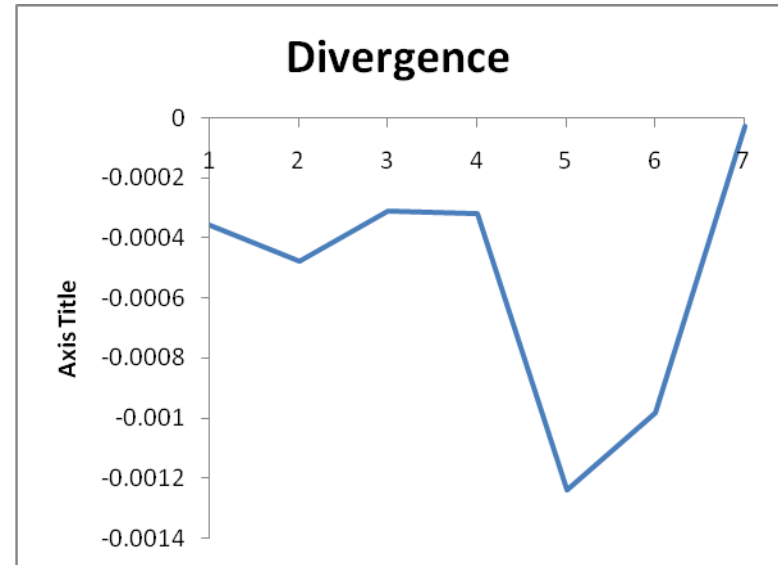
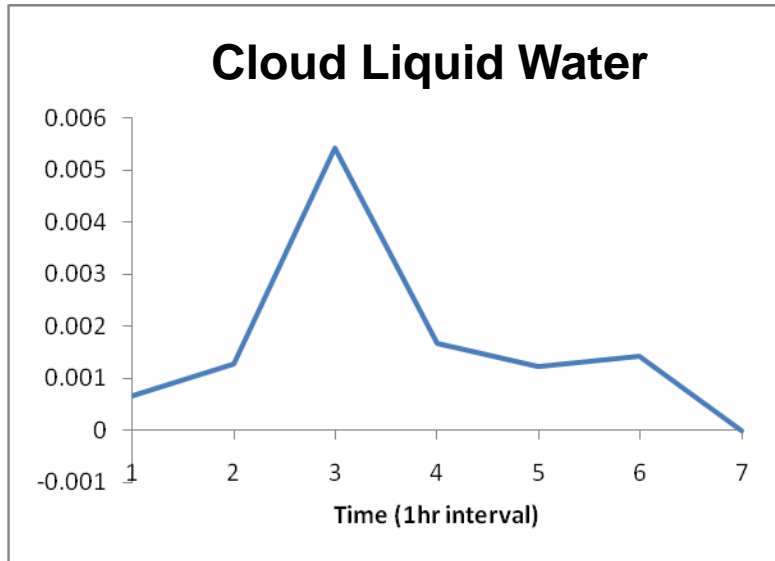


Wind Speed



Initial time: 09z 28 August 2005

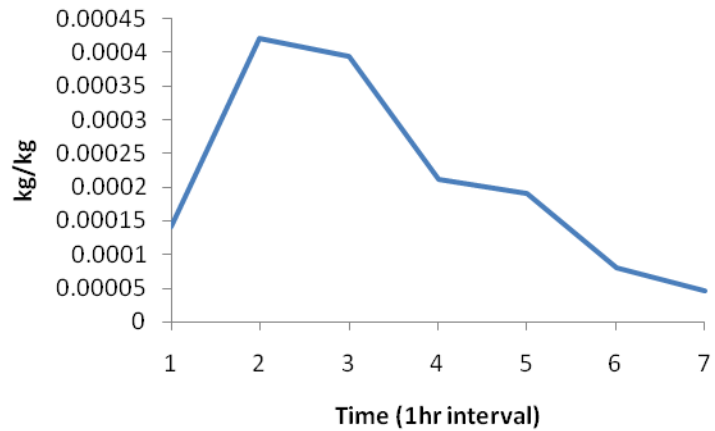
Hourly plots



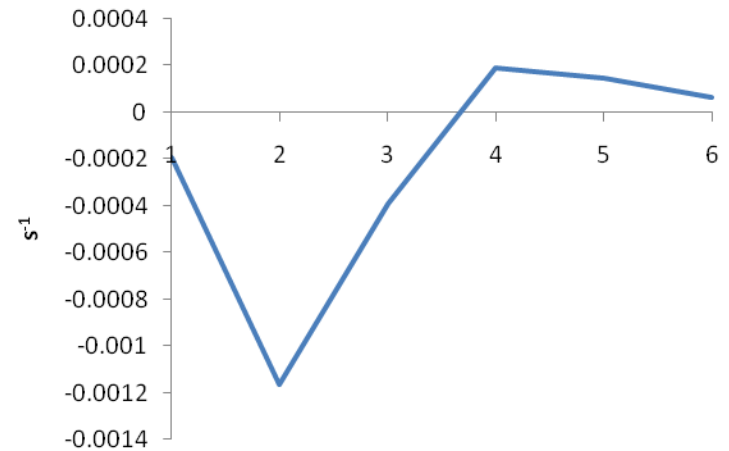
Initial time: 09z 28 August 2005

Hourly plots

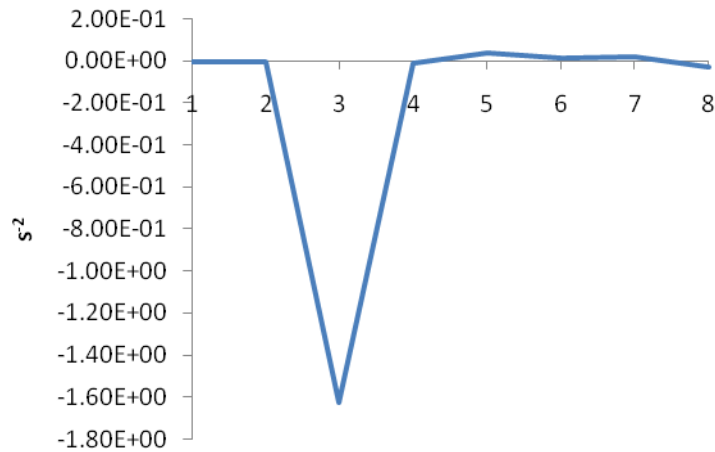
Cloud Liquid Water



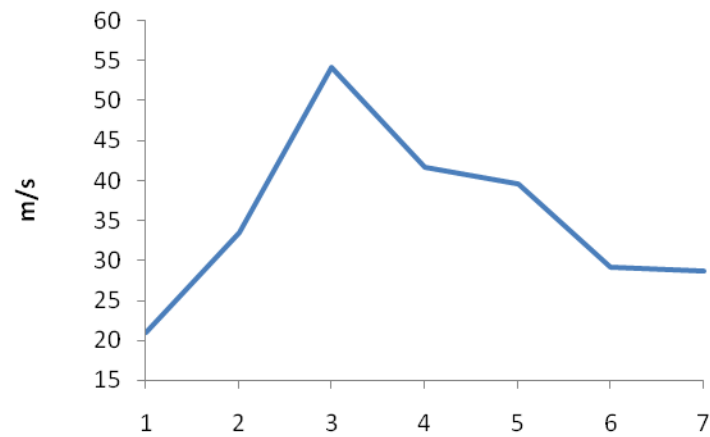
Divergence



Gradient Wind Departure 10⁻⁵

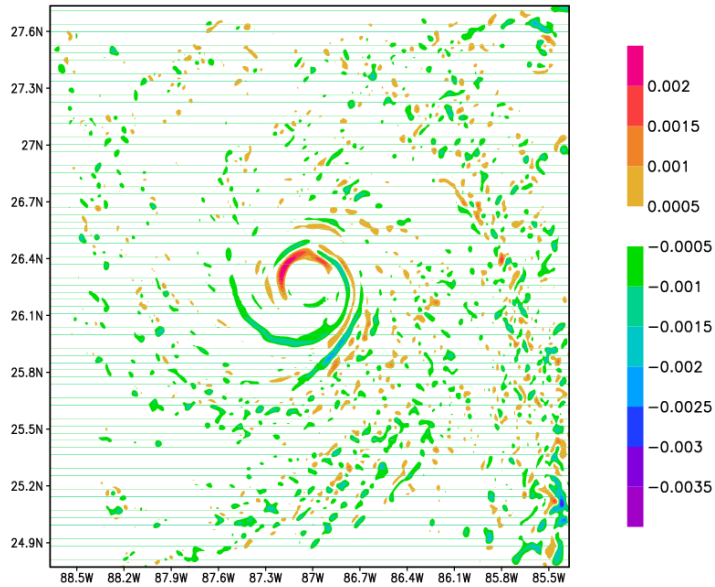


Wind Speed

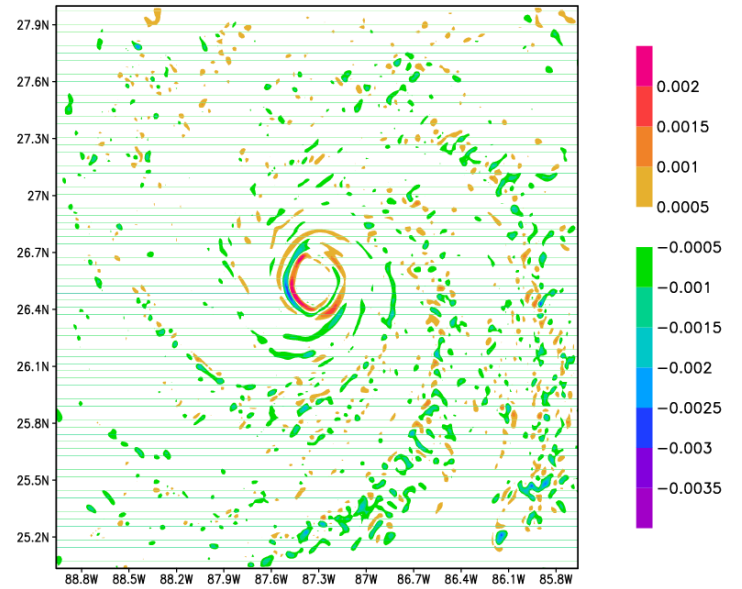


Initial time: 09z 28 August 2005

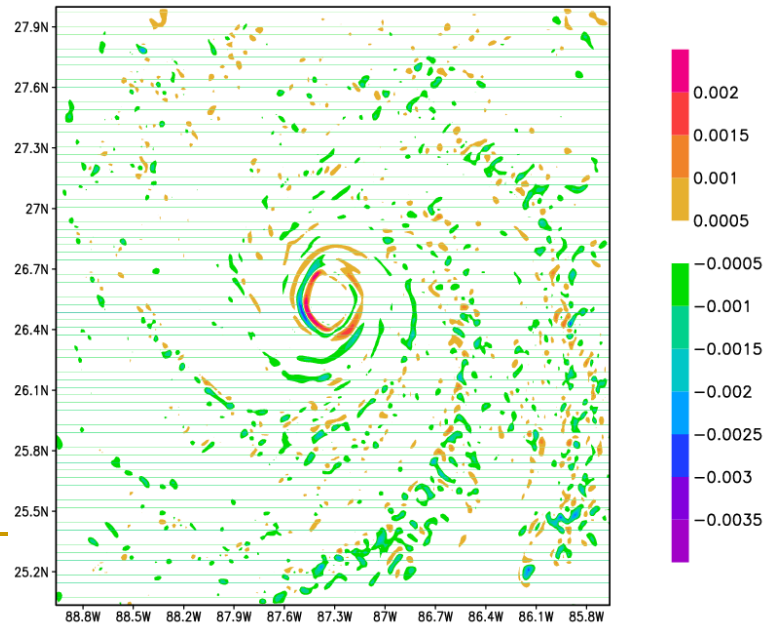
DIVERGENCE (per sec.) 13z 28 Aug 2005



DIVERGENCE (per sec.) 14z 28 Aug 2005

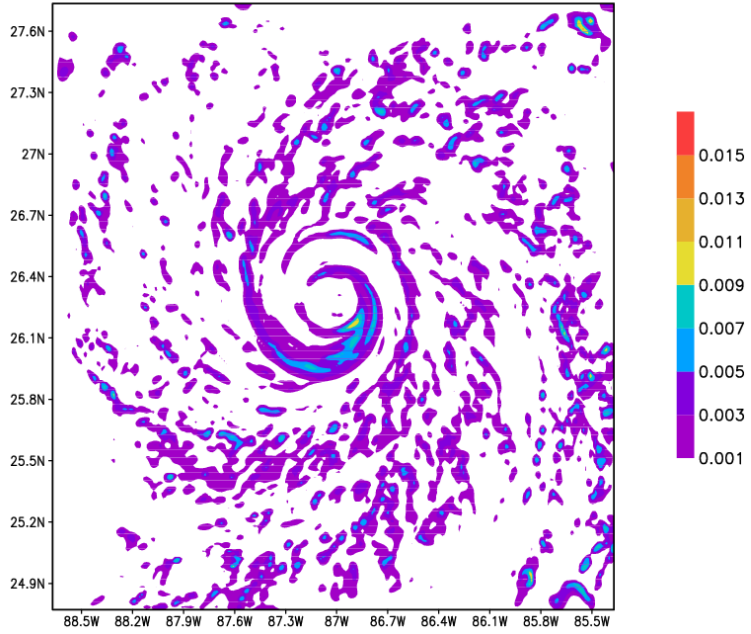


DIVERGENCE (per sec.) 15z 28 Aug 2005

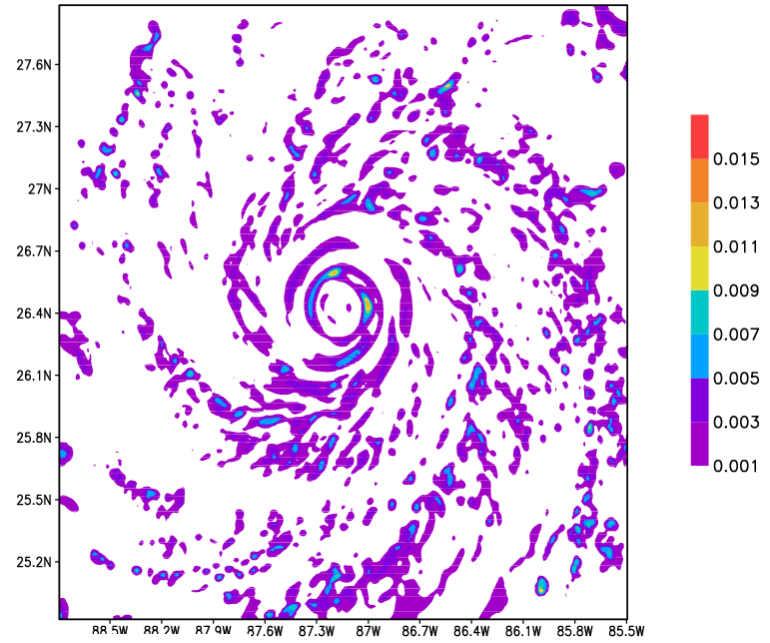


DIVERGENCE

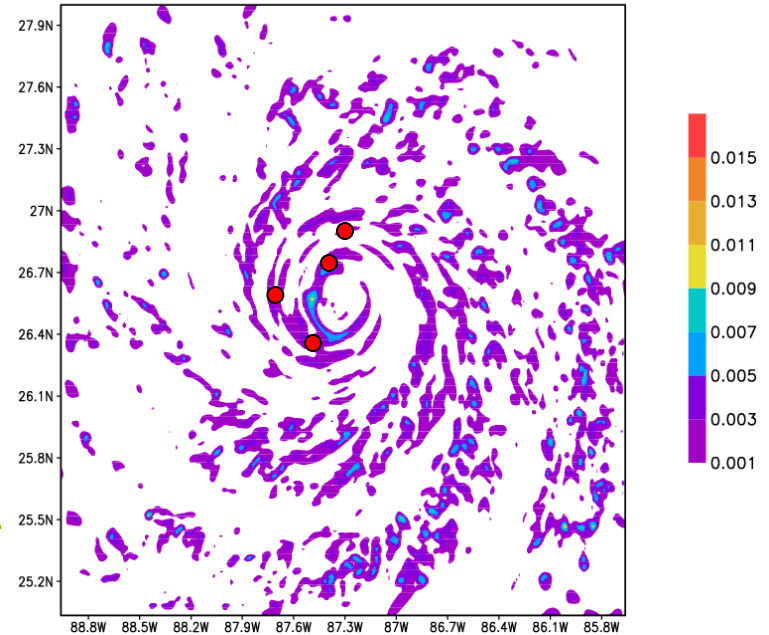
CLOUD LIQUID WATER kg/kg 13z 28 Aug 2005



CLOUD LIQUID WATER kg/kg 14z 28 Aug 2005

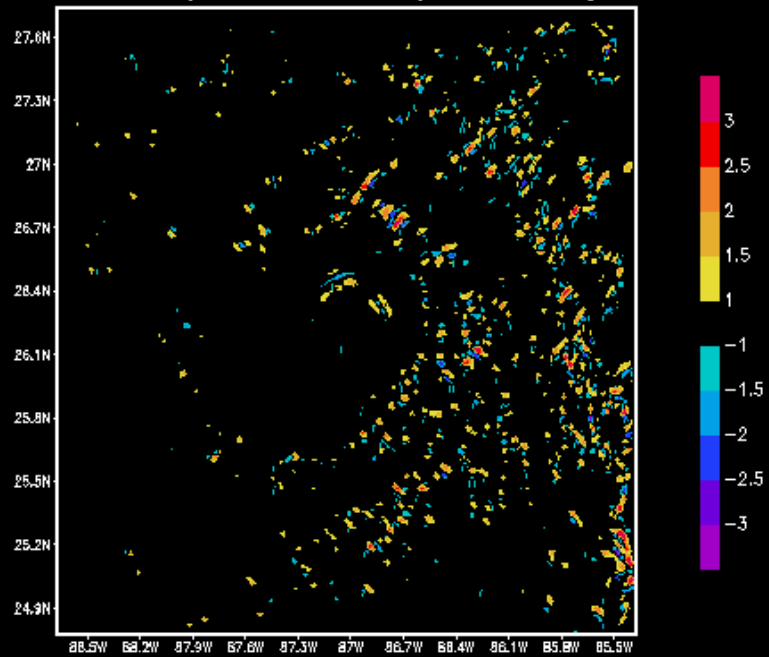


CLOUD LIQUID WATER kg/kg 15z 28 Aug 2005

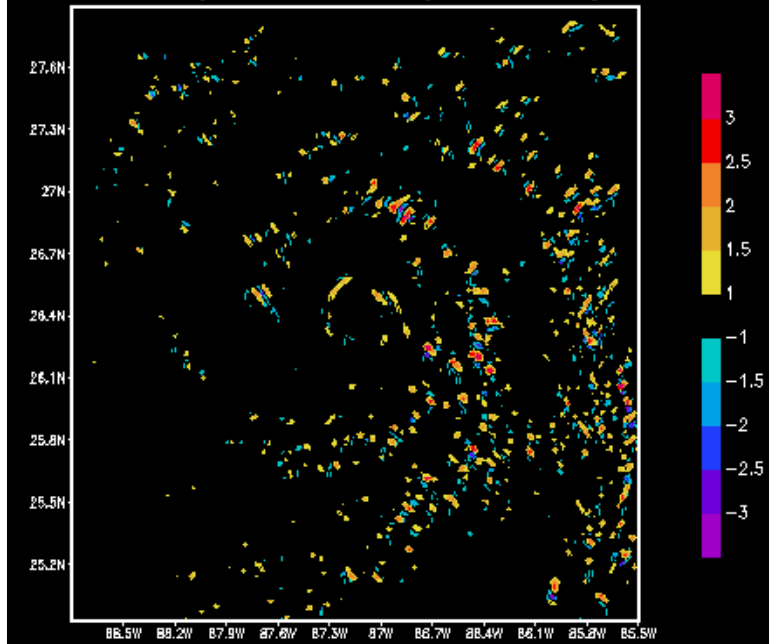


CLOUD LIQUID WATER

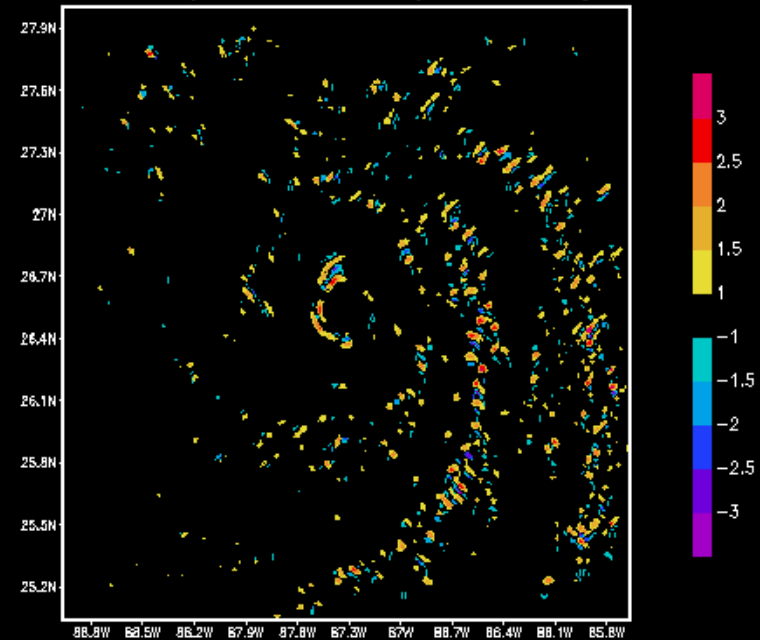
DEPARTURE ($\times 10^{-5} \text{sec}^{**2}$) 13z 28 Aug 2005



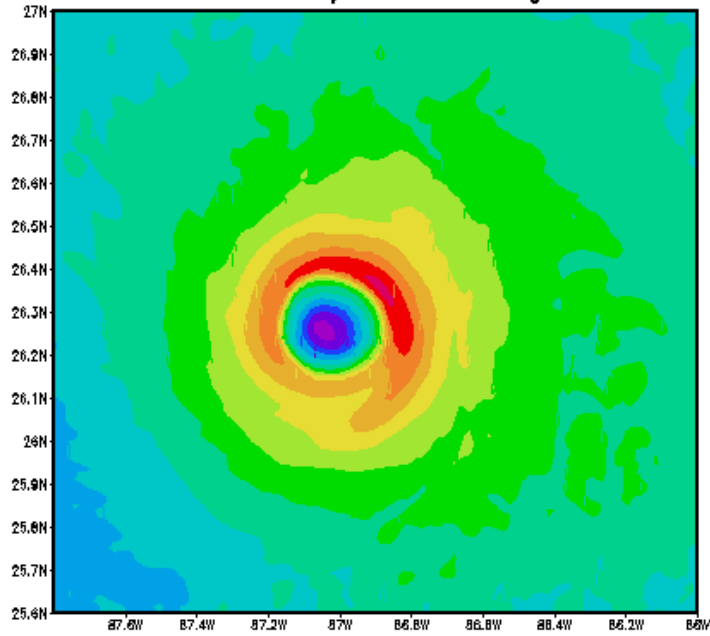
DEPARTURE ($\times 10^{-5} \text{sec}^{**2}$) 14z 28 Aug 2005



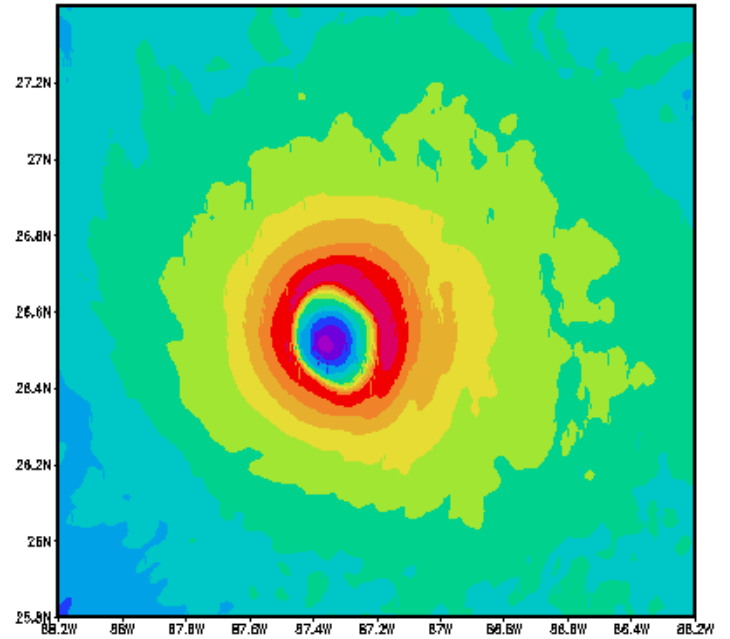
DEPARTURE ($\times 10^{-5} \text{sec}^{**2}$) 15z 28 Aug 2005



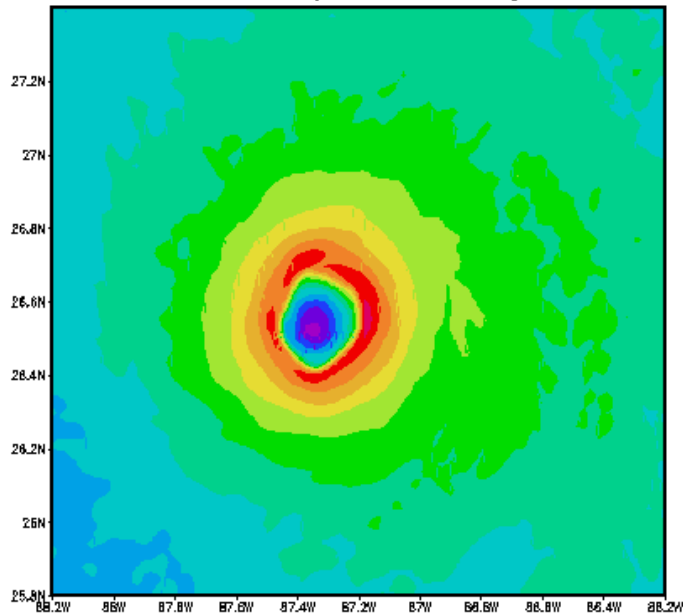
WIND SPEED m/s 13z 28 Aug 2005



WIND SPEED m/s 14z 28 Aug 2005

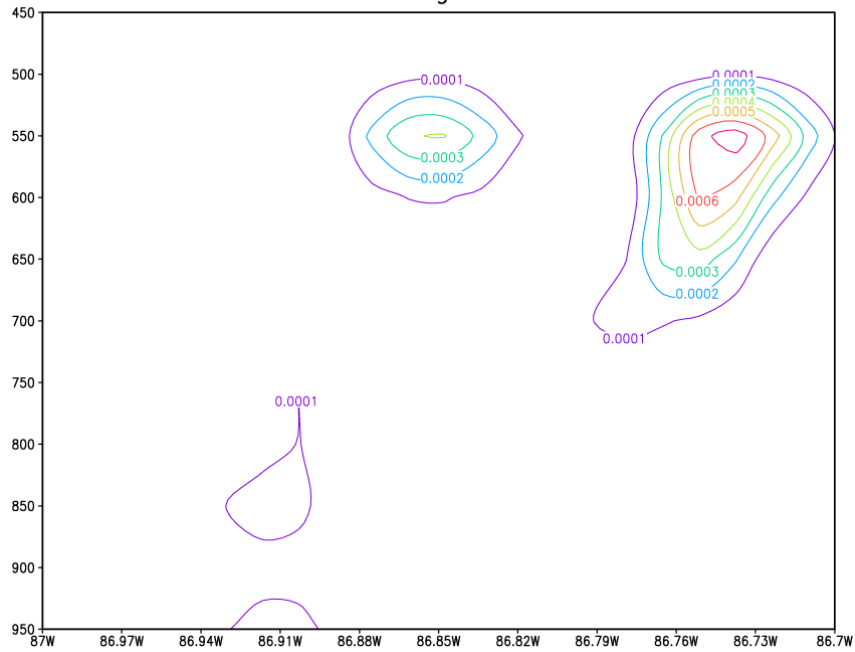


WIND SPEED m/s 15z 28 Aug 2005

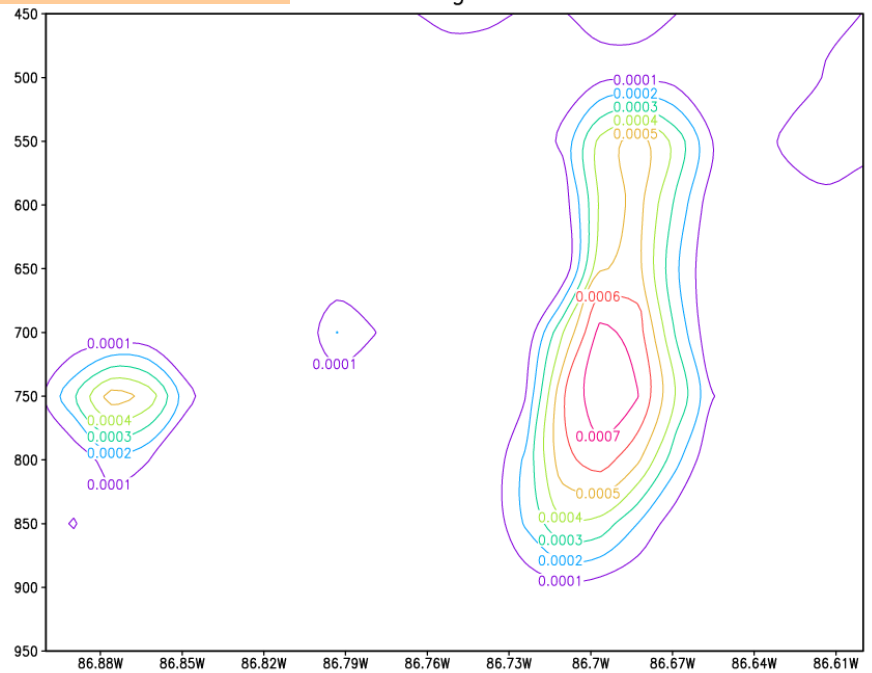


Life cycle of a cloud

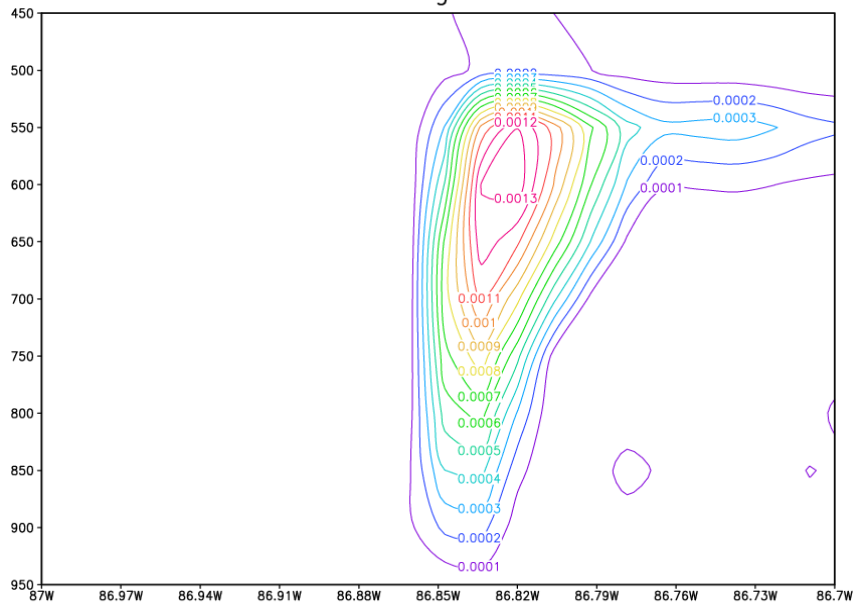
13z 28August 2005



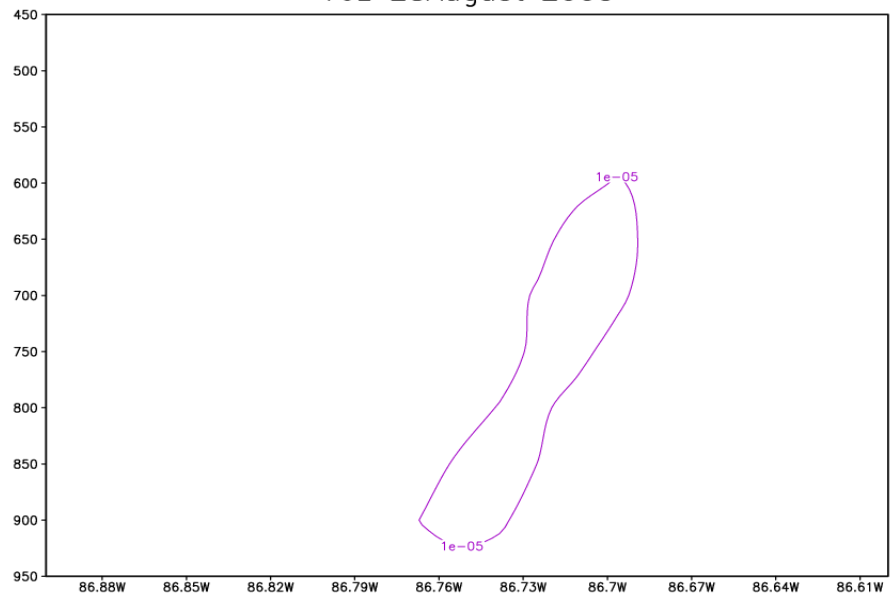
14z 28August 2005



15z 28August 2005

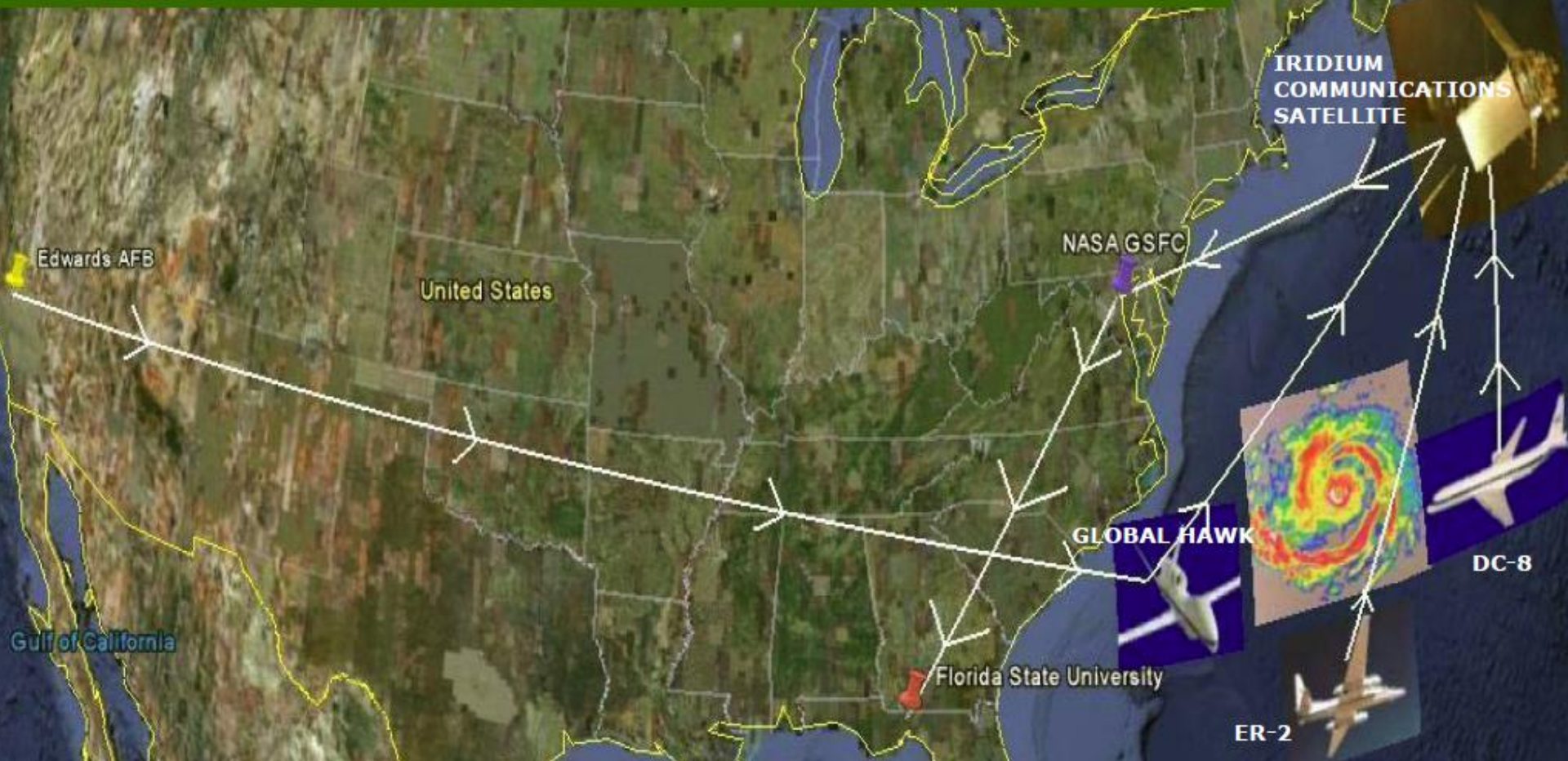


16z 28August 2005



NASA GRIP Experiment- FSU Intensity Forecast Enhancement

- The Global Hawk or the ER-2 leave Dryden/Edwards AFB, or the DC-8 flies into the storm
- The plane observes the storm environment from above
- It transmits data to NASA GSFC via a Communications Satellite
- NASA GSFC transmits the data in near real-time to FSU
- FSU interprets the data to forecast imminent Intensity changes
- FSU sends the results to Mission Scientists



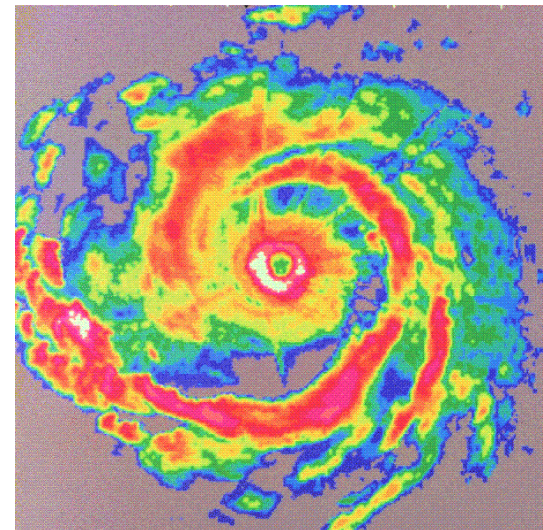
Observations Required

- Radar Reflectivity
- 3-Dimensional Winds
- Pressure Altitude

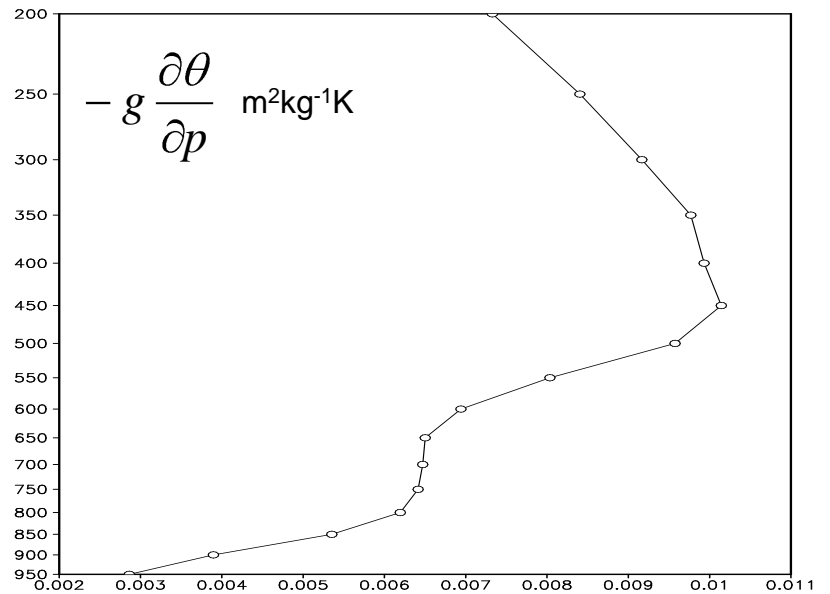
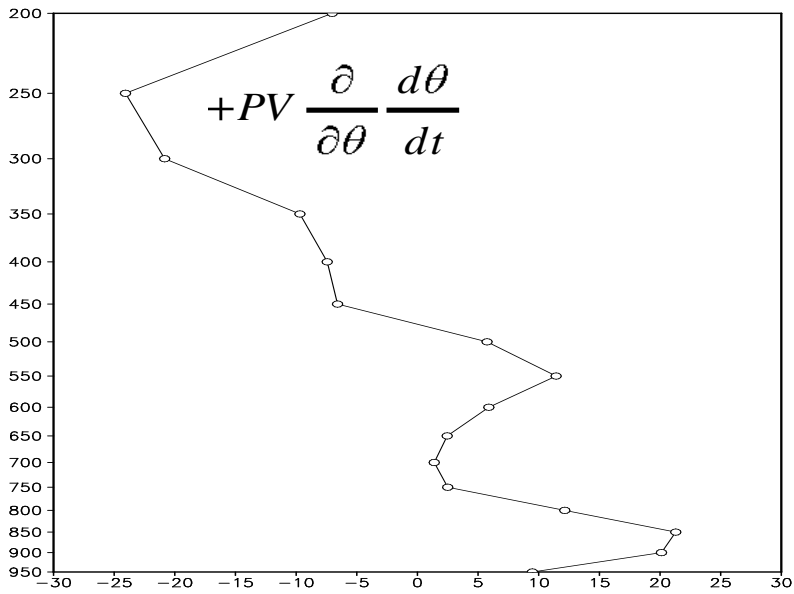
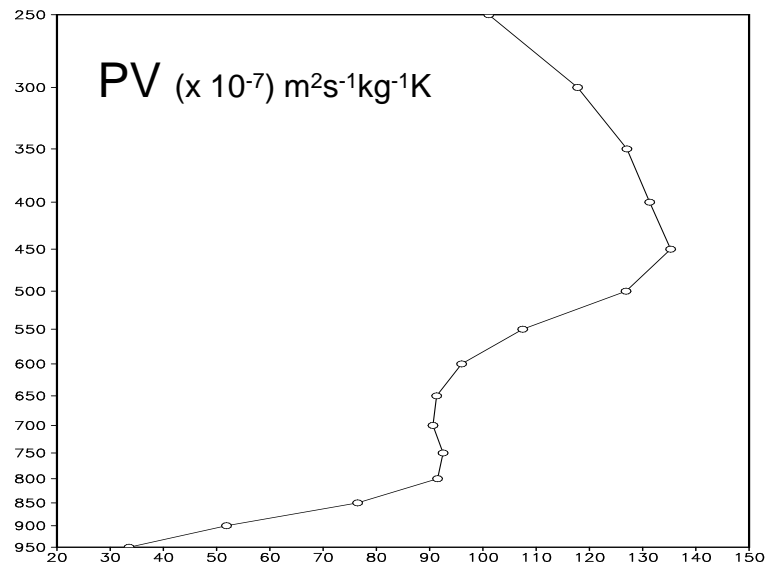
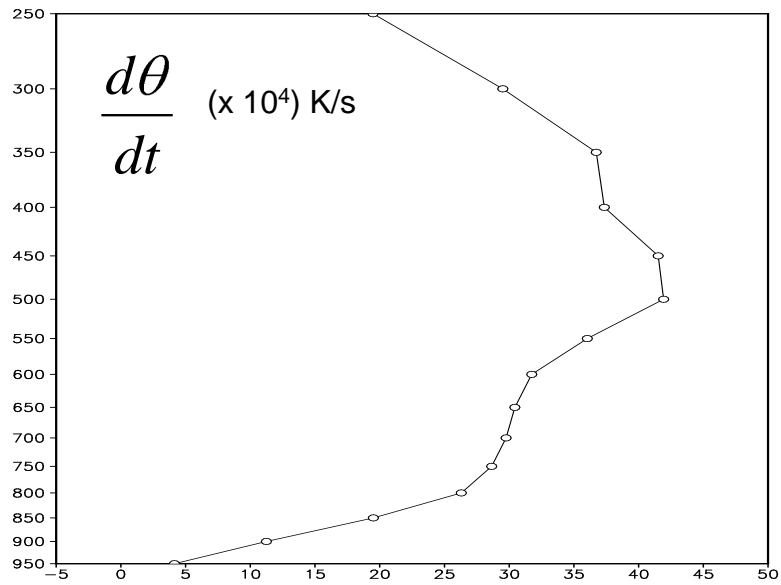


NASA Dryden Flight Research Center Photo Collection
<http://www.dfrc.nasa.gov/Gallery/Photo/index.html>
NASA Photo: ED08-0309-24 Date: December 11, 2008 Photo By: Tony Landis

This and a companion Global Hawk unmanned aircraft are used by NASA for Earth science missions and by Northrop Grumman for developmental testing.

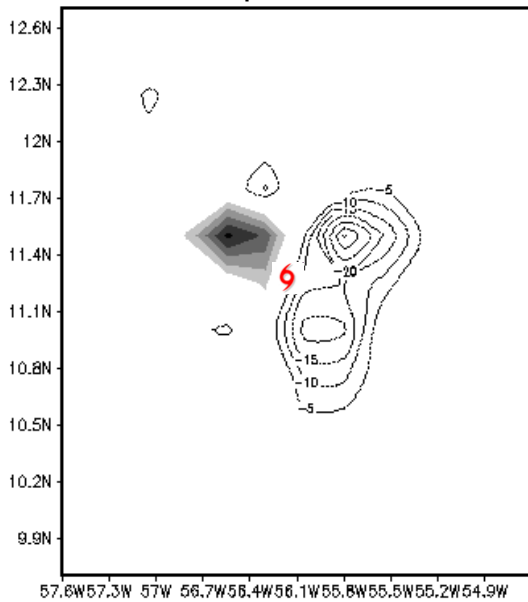


**Vertical differential of heating in the
complete potential vorticity equation,
a diagnostic tool**

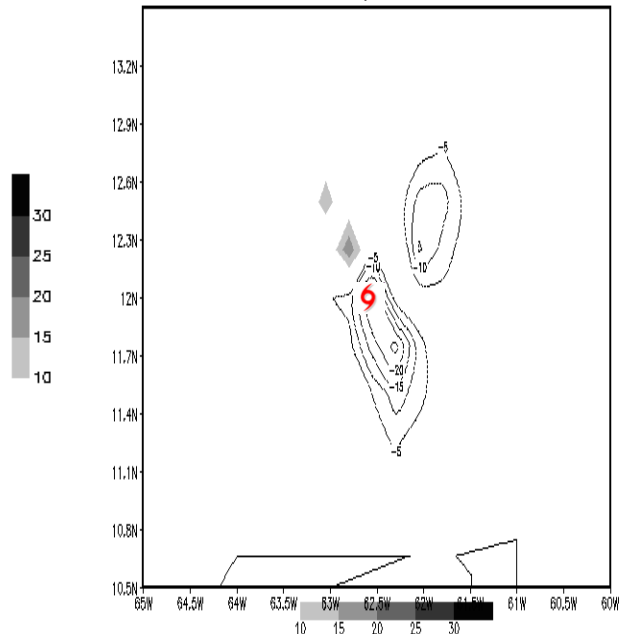


Hurricane IVAN, 11 September 2004, 12z. 10^{-7}) $\text{m}^2\text{s}^{-1}\text{kg}^{-1}\text{K}$

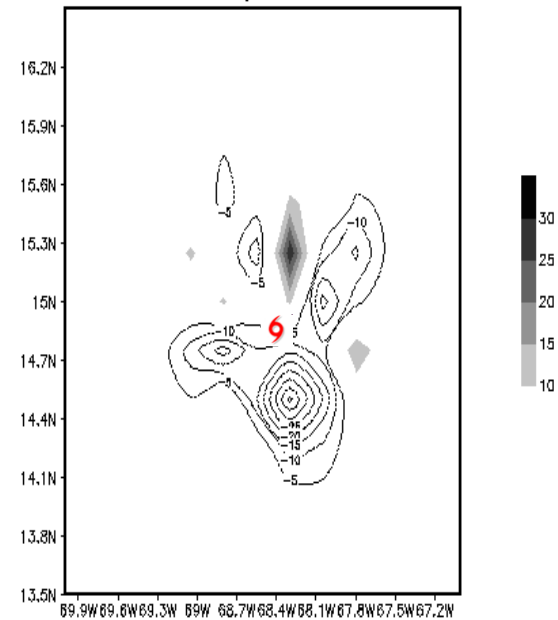
7 September 00z



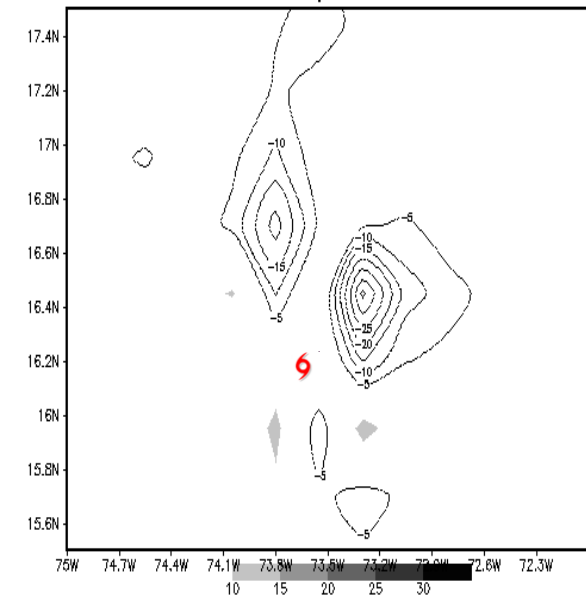
8 September



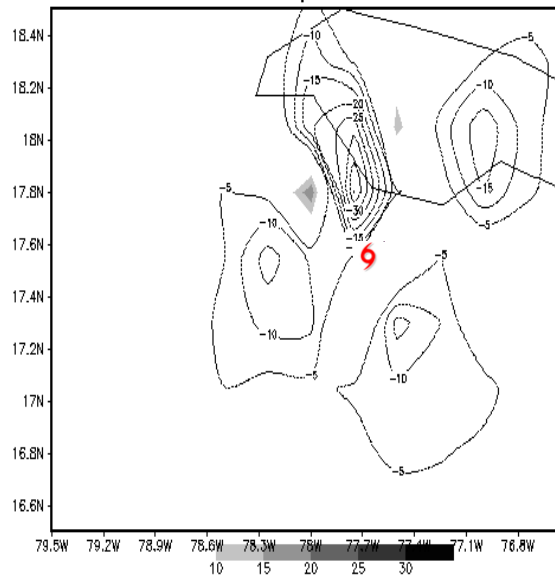
9 September



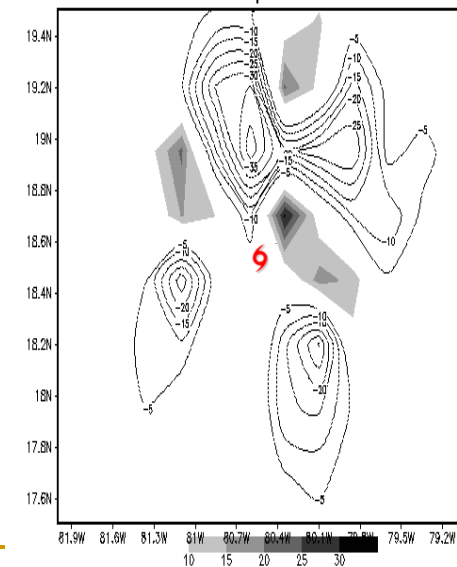
10 September



11 September

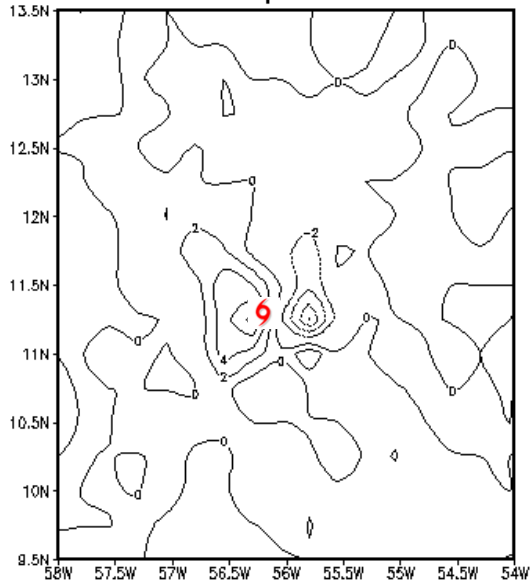


12 September

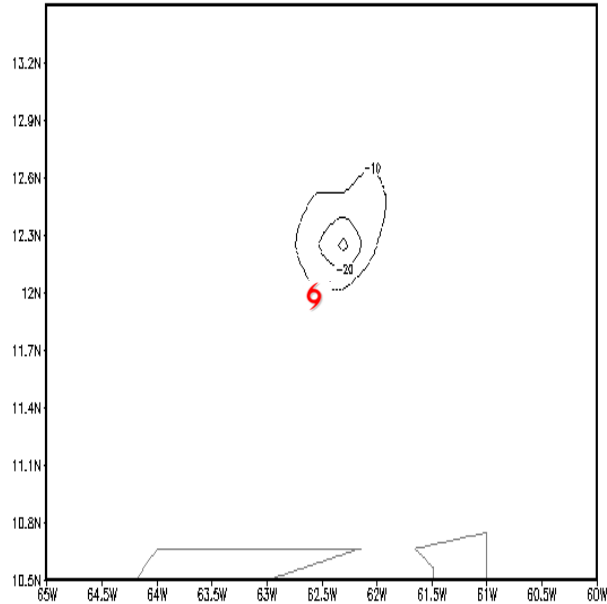


Contour plots of horizontal advection (x 10⁻¹⁰ Kg⁻¹m⁻²s⁻²K) for hurricane IVAN 7 September through 12 September 2005 at 00z.

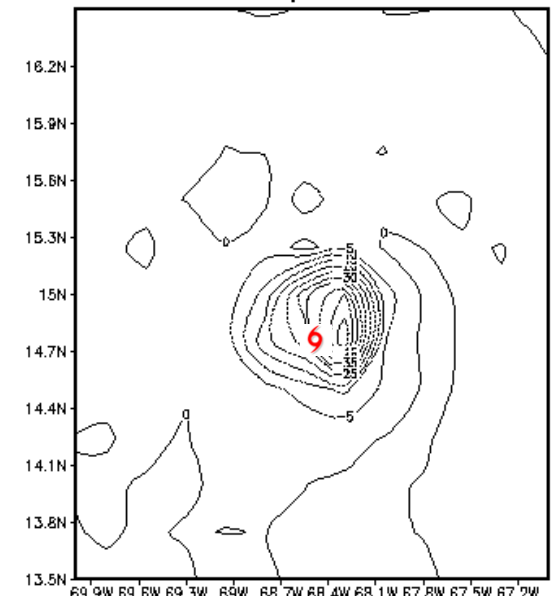
7 September



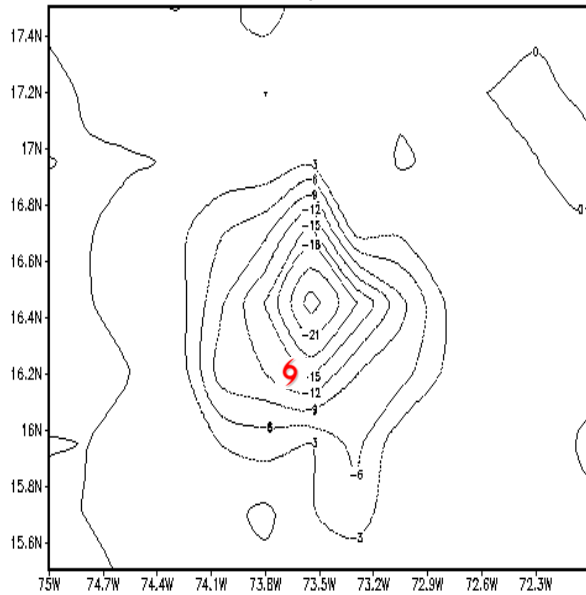
8 September



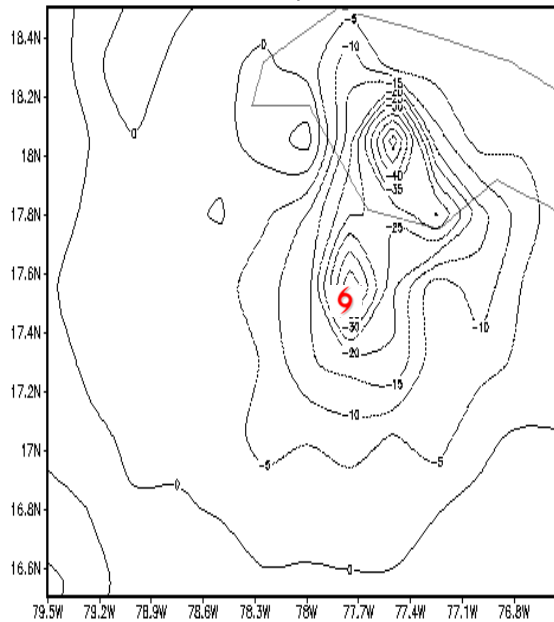
9 September



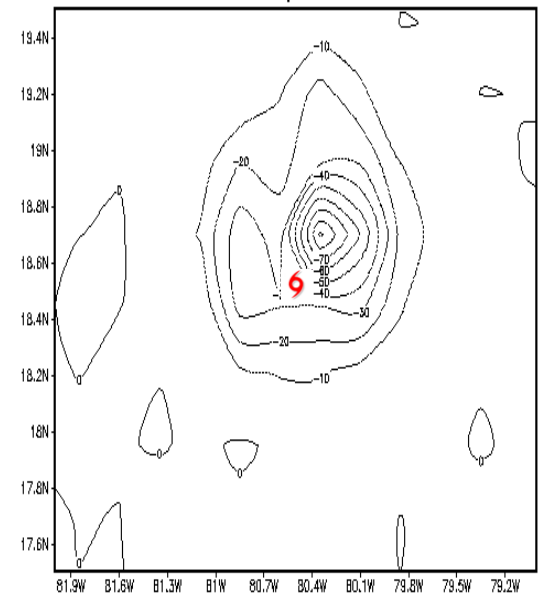
10 September



11 September

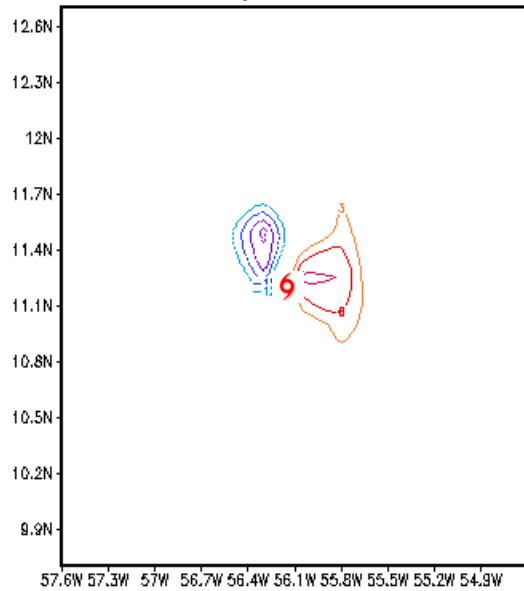


12 September

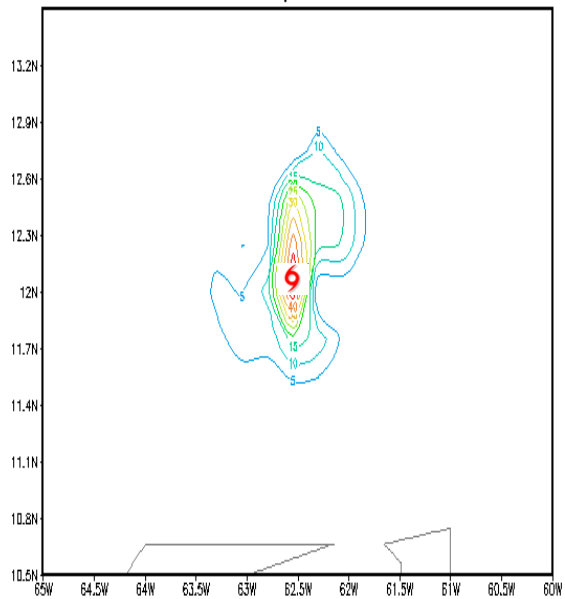


Contour plots of vertical advection of $PV(x 10^{-10} \text{ Kg}^{-1}\text{m}^2\text{s}^{-2}\text{K})$ for hurricane IVAN, 7 September through 12 September 2005 at 00z.

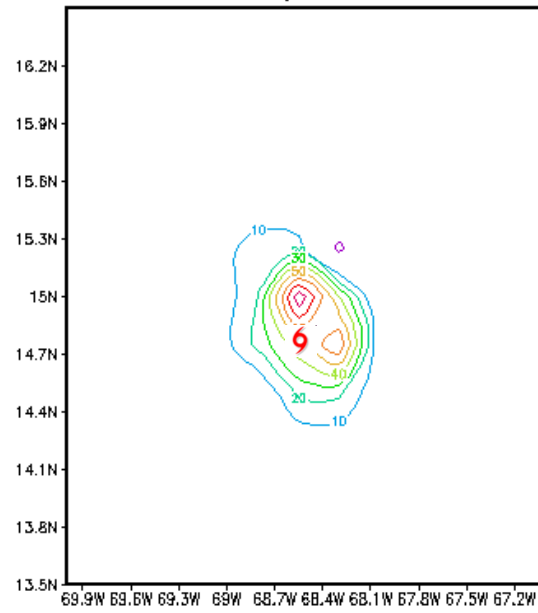
7 September 00z



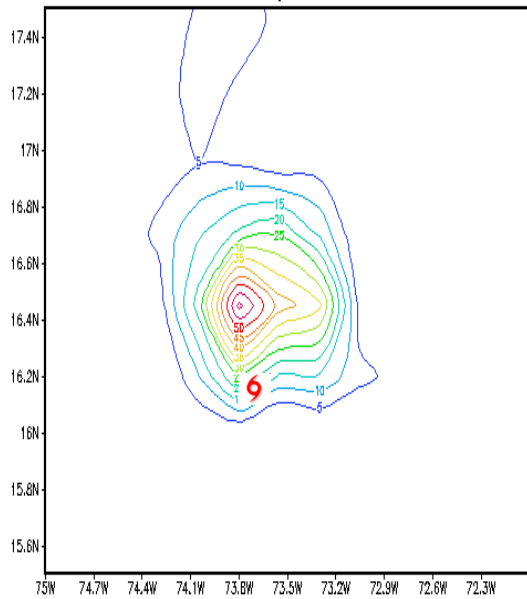
8 September



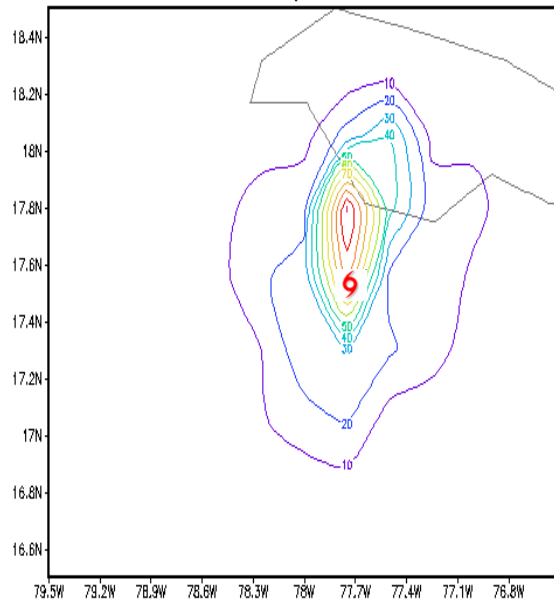
9 September



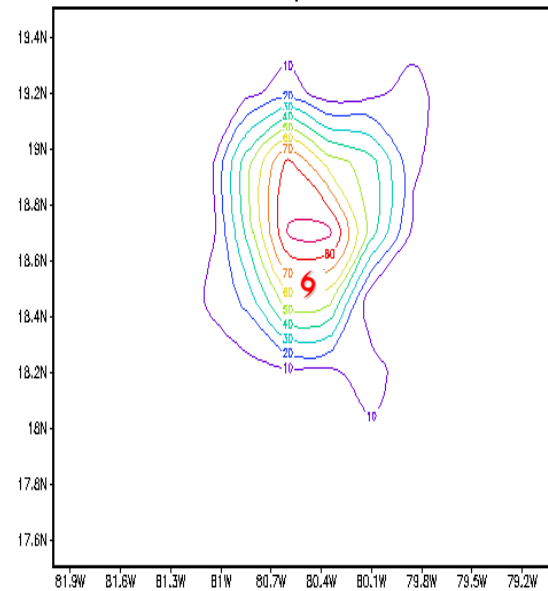
10 September



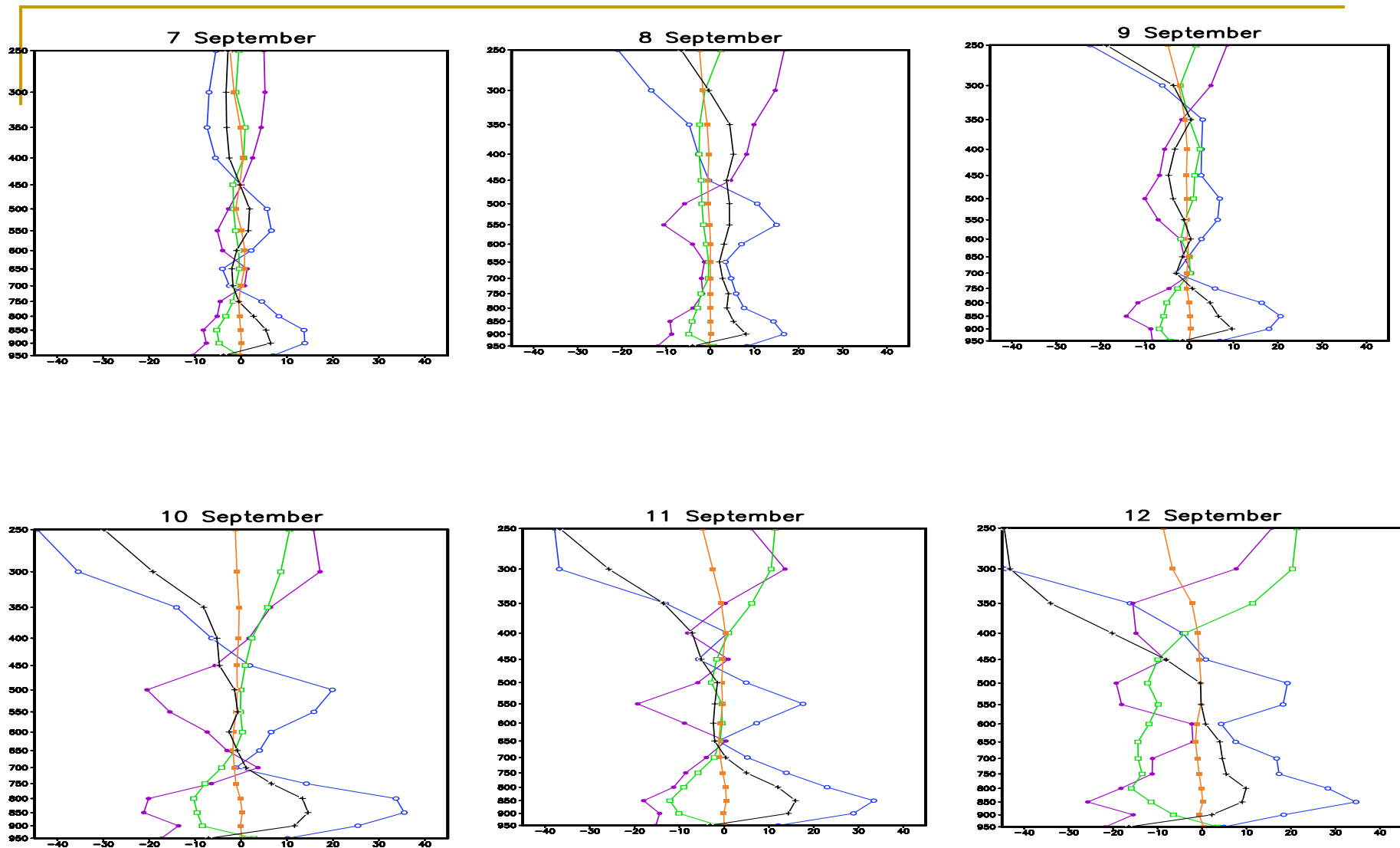
11 September



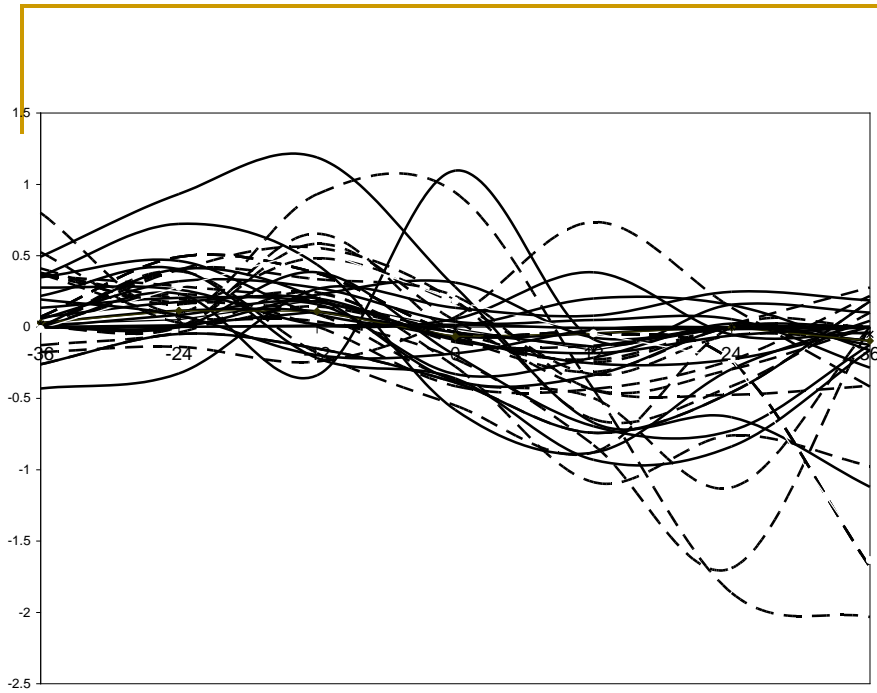
12 September



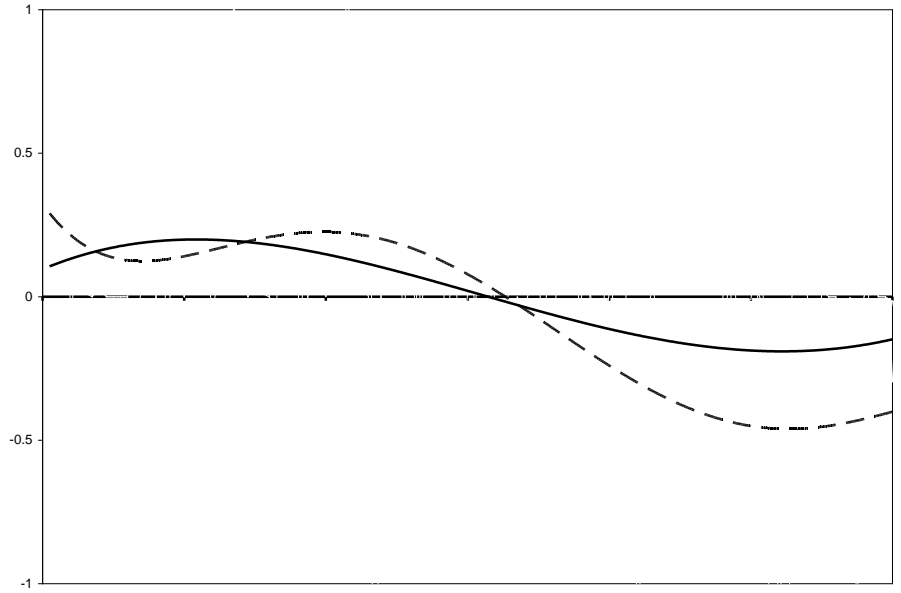
Contour plots of vertical differential of heating (x 10⁻¹⁰ Kg⁻¹m²s⁻²K) for hurricane IVAN, 7 September through 12 September 2005 at 00z.



Vertical distribution of the Horizontal Advection (green), Vertical Advection (purple), Vertical Differential of heating (Blue), Horizontal Differential of heating (orange) for the Potential Vorticity during the intensifying stage of Hurricane IVAN, through 7 to 12, 2005 at 00z. The total diabatic heating is shown in black. Units are $\times 10^{-10} \text{ Kg}^{-1}\text{m}^2\text{s}^{-2}\text{K}$



The time rate of change of vertical differential of heating (black solid line) and that for intensity (black dashed line) for the individual hurricanes for the years 2004-2006. The abscissa is a relative time scale and ordinate denotes a scale for the time rate of change.



Least square fit of the above two curves.

We shall next show how the rapid increase of divergent kinetic energy gets transformed into rotational kinetic energy leading to a stronger storm

Hurricane Katrina

8/28/05 12Z – 8/29/05 15Z

Every 3 hours

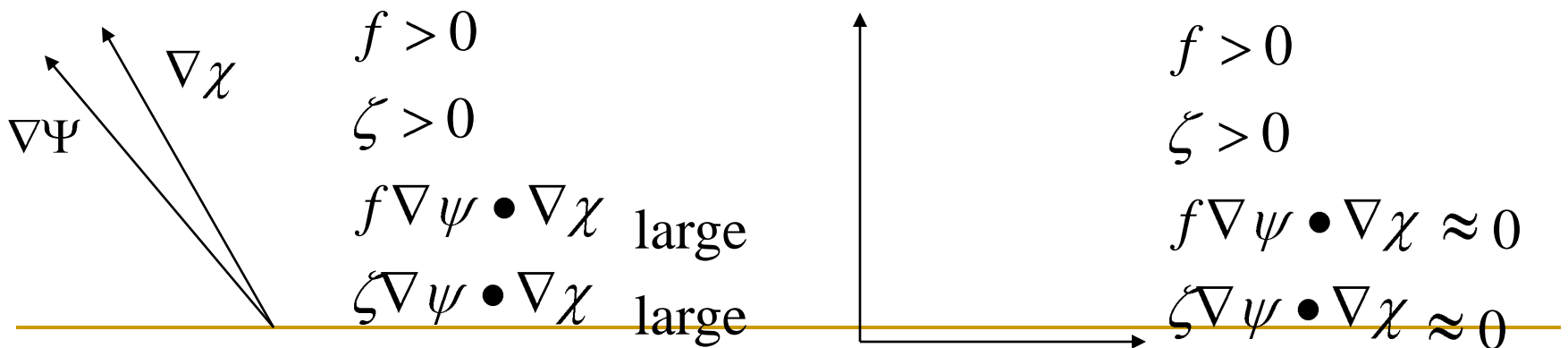
And several other hurricanes during 2004
- 2006

Equations for Psi-Chi Interactions

$$\frac{\partial}{\partial t} \overline{K_\psi} = B_\psi + \overbrace{f \nabla \psi \cdot \nabla \chi}^{\text{Term 1}} + \overbrace{\nabla^2 \psi \nabla \psi \cdot \nabla \chi}^{\text{Term 2}} + \overbrace{\nabla^2 \chi (\nabla \psi)^2 / 2}^{\text{Term 3}} + \overbrace{\omega J \left(\psi, \frac{\partial \chi}{\partial p} \right)}^{\text{Term 4}} + F_\psi$$

The important terms are Term 1 and Term 2

$$\nabla \psi \cdot \nabla \chi$$

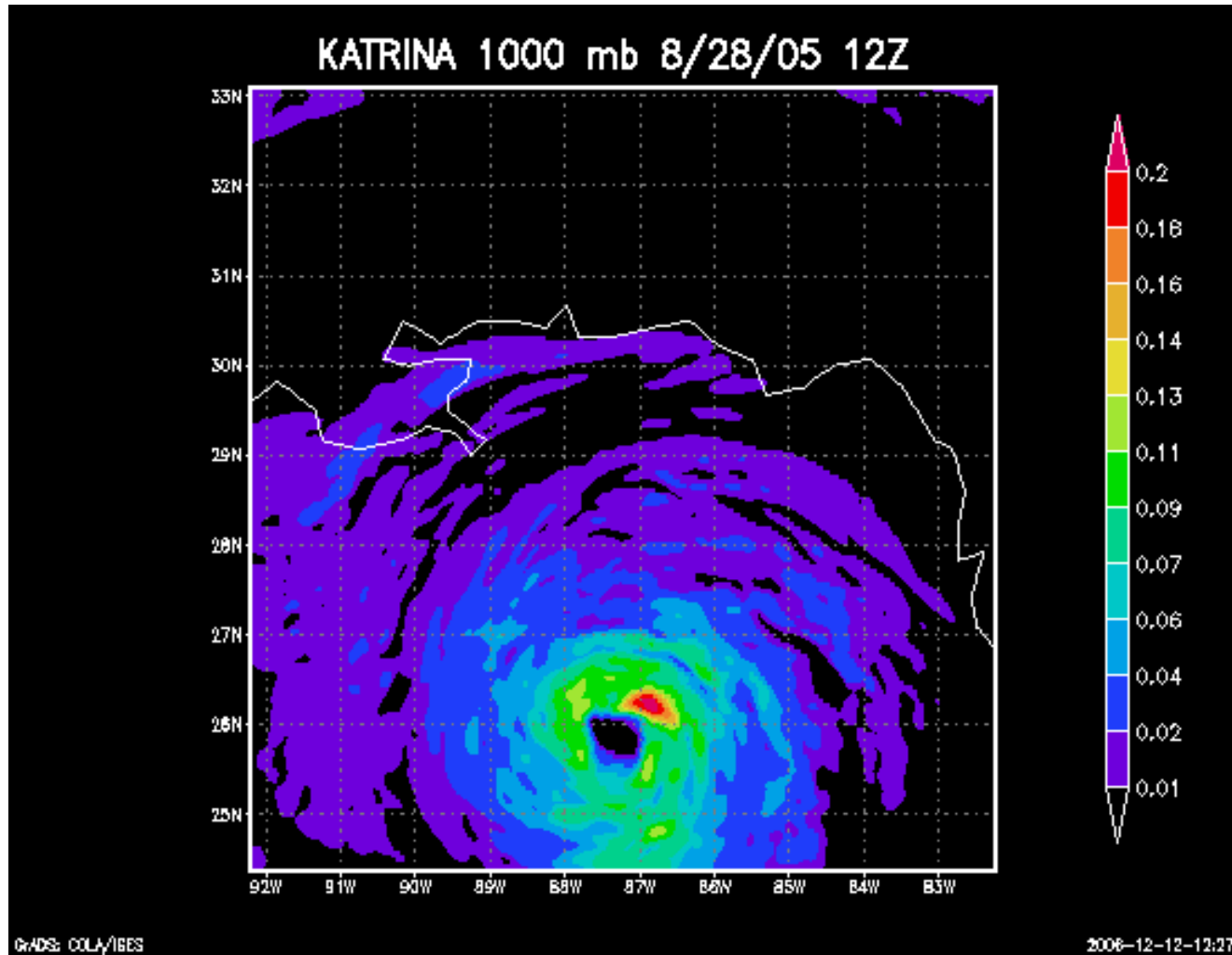


$$\frac{\partial}{\partial t} \overline{\overline{K_\psi}} = B_\psi + \overline{\overline{\text{Term 1}}} + \overline{\overline{\text{Term 2}}} + \overline{\overline{\text{Term 3}}} + \overline{\overline{\text{Term 4}}} + F_\psi$$

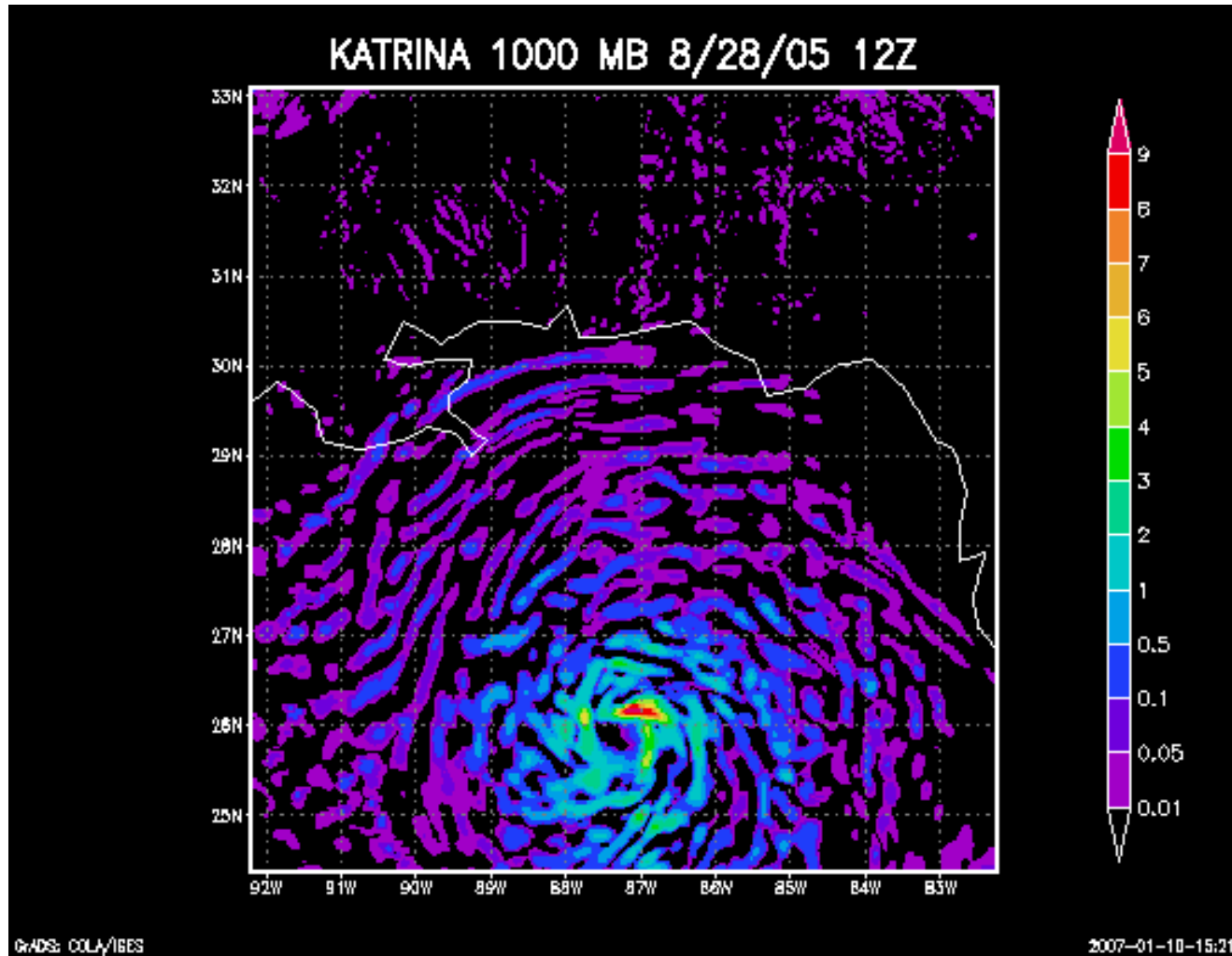
$$\frac{\partial}{\partial t} \overline{\overline{K_\psi}} = B_\psi + \overline{\overline{f \nabla \psi \cdot \nabla \chi}} + \overline{\overline{\nabla^2 \psi \nabla \psi \cdot \nabla \chi}} + \overline{\overline{\nabla^2 \chi (\nabla \psi)^2 / 2}} + \overline{\overline{\omega J \left(\psi, \frac{\partial \chi}{\partial p} \right)}} + F_\psi$$

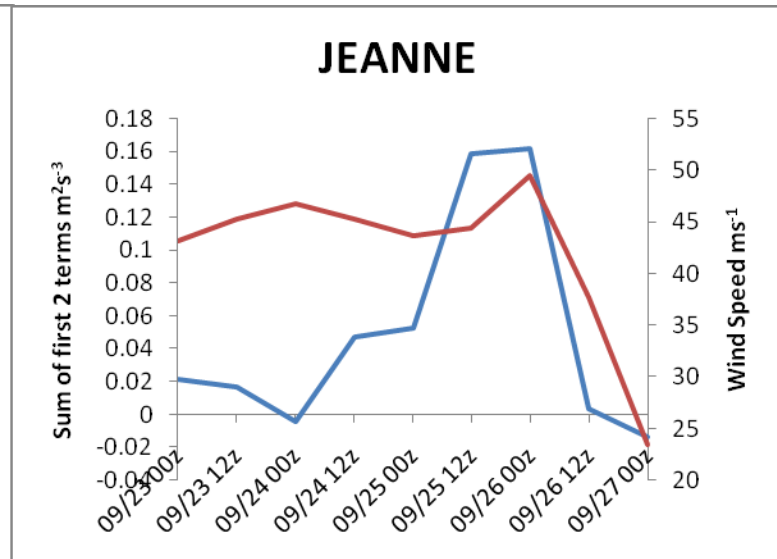
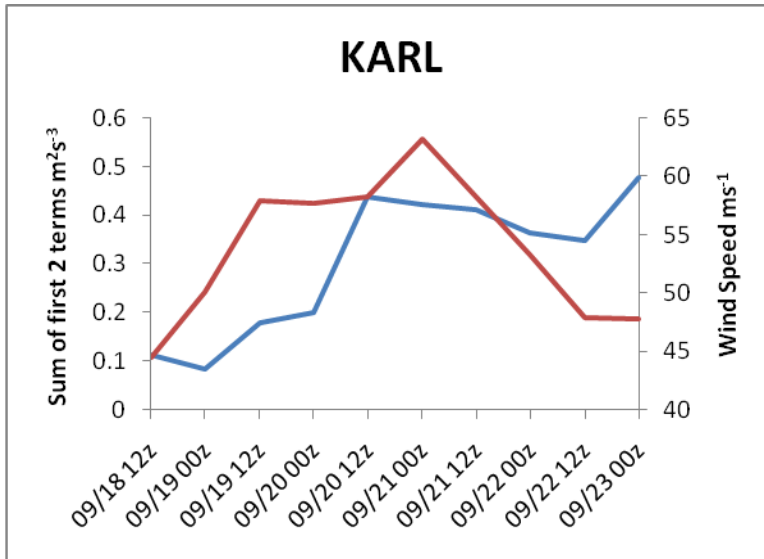
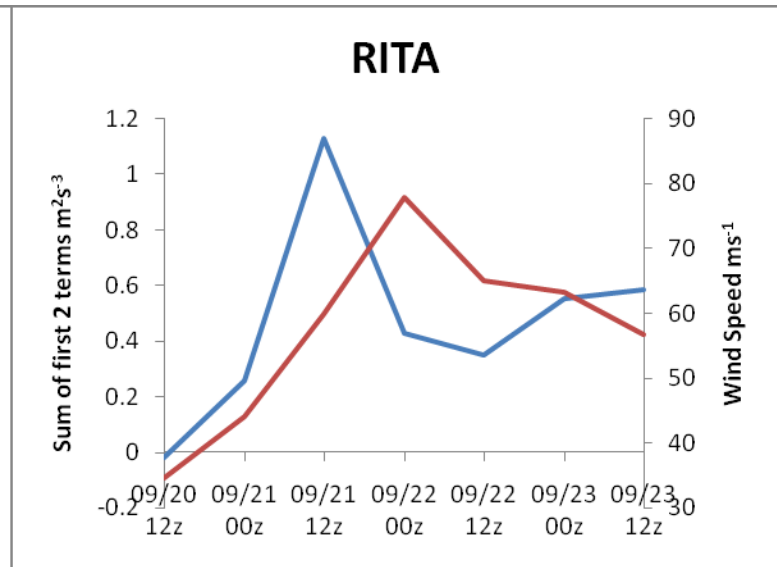
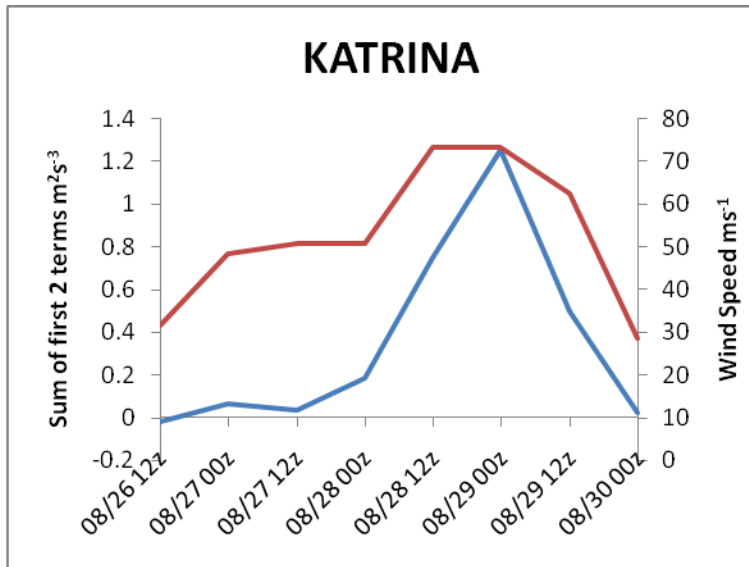
There are four terms in the psi-chi interactions, these measure the rate of transfer of divergent kinetic energy into rotational kinetic energy. As deep convection flares up divergent motions amplify and a strong conversion of divergent to rotational motion follows, continually imbalanced flows grow and a rapid intensification of hurricane winds follow.

$$f(\nabla\psi \cdot \nabla\chi) \quad (\text{m}^2\text{s}^{-3})$$

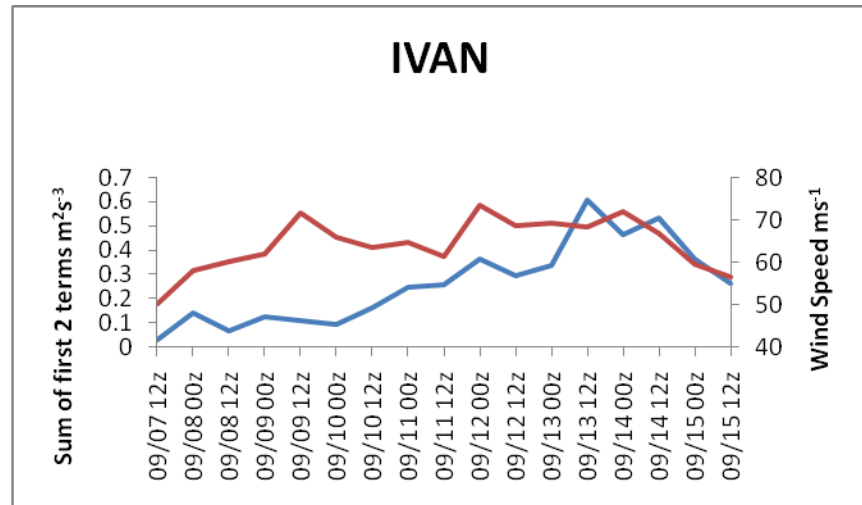
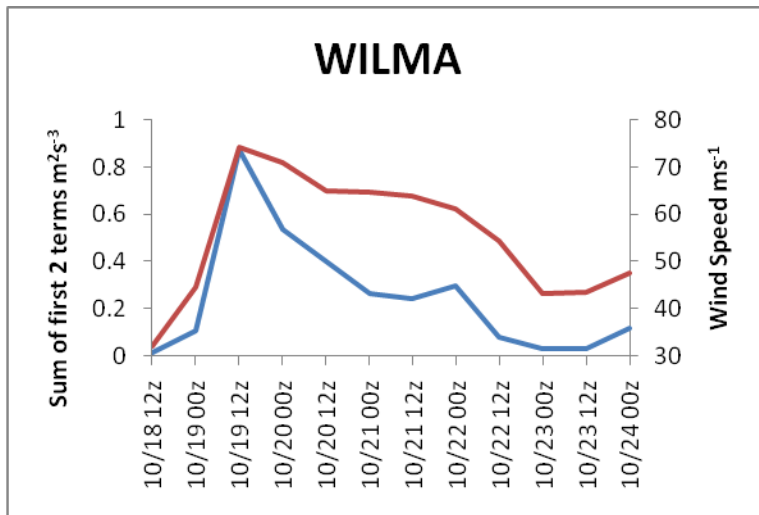
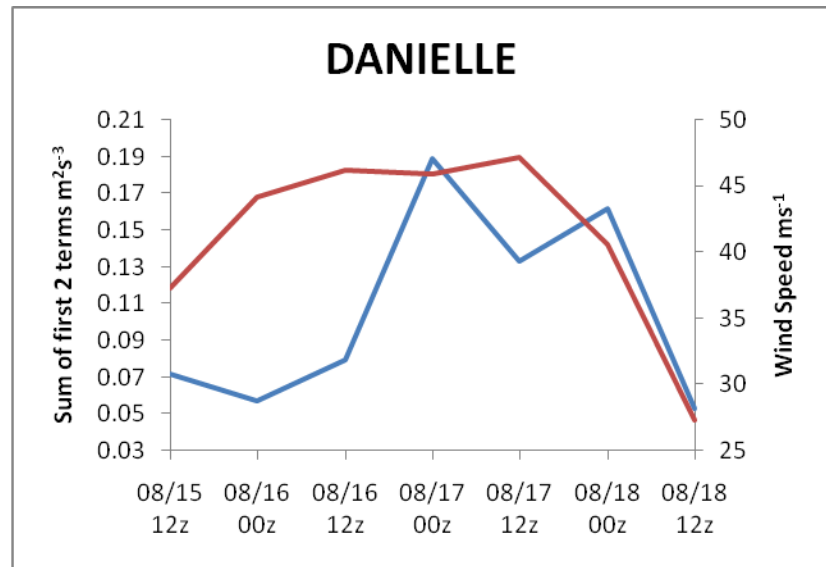
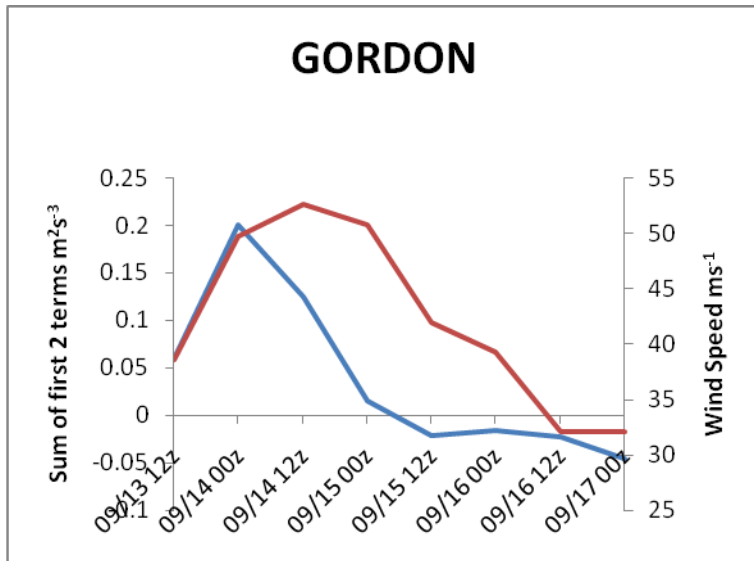


$$\nabla^2 \psi (\nabla \psi \bullet \nabla \chi) \quad (\text{m}^2\text{s}^{-3})$$





These results are averages over 100 km radius at 850 hPa
 Red line is intensity; Blue line is the sum of the first two terms



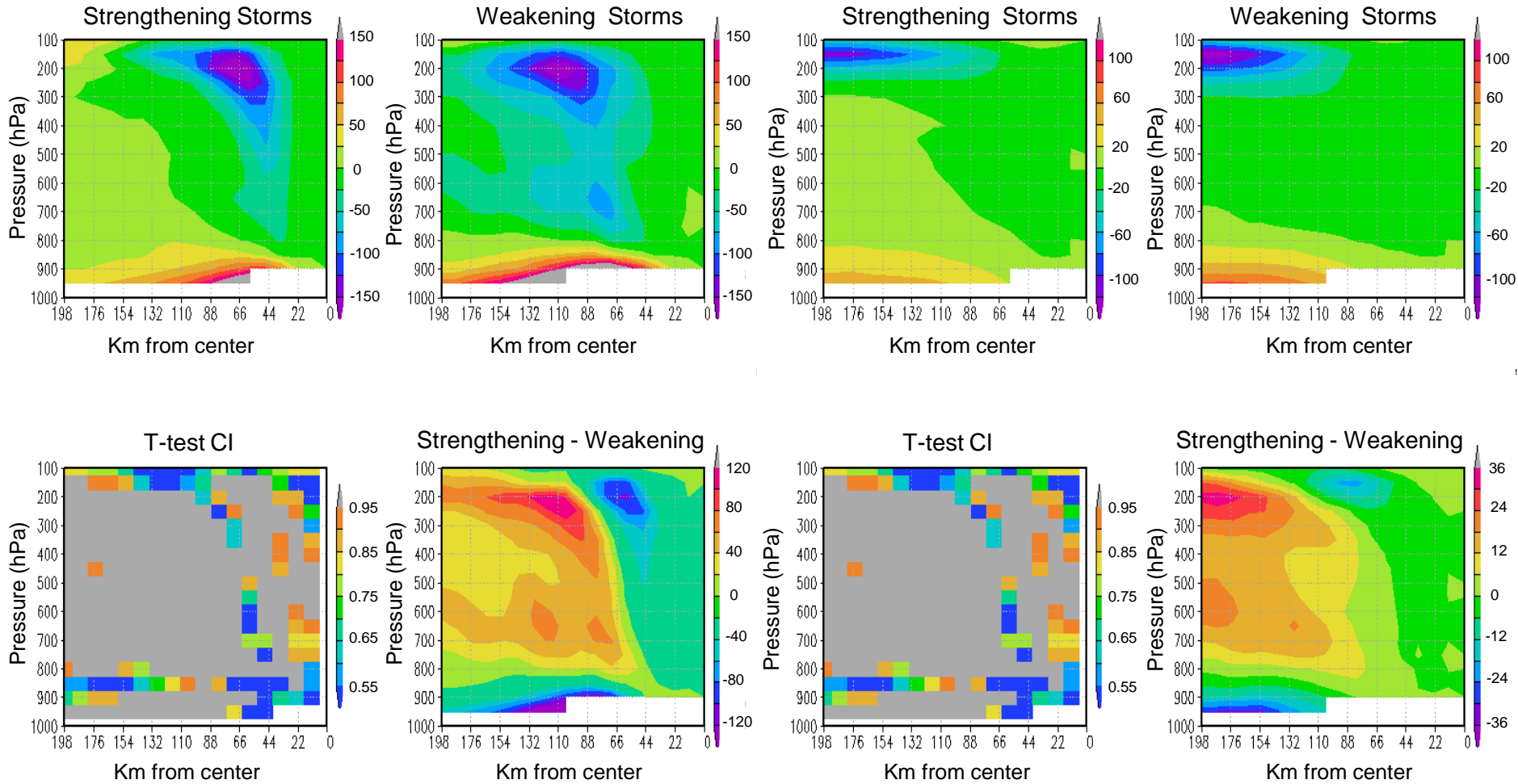
These results are averages over 100 km radius at 850 hPa
 Red line is intensity; Blue line is the sum of the first two terms

Storm-Relative Eulerian Absolute Angular Momentum Tendencies
in Atlantic Tropical Cyclones

Cross-section Composites of Horizontal Advection in Storms Category 2 and Higher

$$-\vec{V} \cdot \nabla (u_{\theta}) \text{ m}^2 \text{ s}^{-2}$$

$$-\vec{V} \cdot \nabla \left(\frac{f r^2}{2} \right) \text{ m}^2 \text{ s}^{-2}$$



Conclusions and future work contd..

- Future work on mesoscale modeling during the hurricane season of 2009 will include the following models: HWRF(EMC), HWRF-X(HRD), WRF/ARW (NCAR), COAMPS (NRL),GFDL(NOAA), MM5 (FSU), WRF(FSU).

- Rapid intensity changes over regions of cloud burst seem to go with growth of lower tropospheric convergence, growth of departure from balance laws and supergradient winds. This scenario appears to carry a lead time of only an hour (roughly) between the time of the cloud burst and the generation of supergradient winds. Local wind maxima thus generated are seen to sweep azimuthal distances of the order of 50 km in a matter of an hour.

- There are a host of dynamical parameters that deserve to be examined over the inner core of hurricanes for the post processing of model output. We have noted the following to be important, in relating to hurricane intensity changes:
 - vertical differential of heating of the complete pv equation;
 - energy transfer from the divergent to the rotational kinetic energy;
 - advection of earths and relative angular momentum in storm relative coordinates;
 - contributions of the kinematic transfer of shear vorticity to the curvature vorticity.

Future work

- **In order to further exploit these findings, we have designed a unified index that is derived from multiple regression for the intensity change using a least square minimization principle that includes all of the above parameters. Such an index would be most useful for real time operations. It should be stated that the parameters we have selected are not all-inclusive; such a list must be expanded and tested as newer ideas develop. The post spin-up data sets are ideally suited for developing single or combined indices for studying the rapid intensification or weakening. A combined index is described by the following multiple regression equation.**

$$\text{Intensity Change} = \sum \alpha_i \text{PR}_i \quad (1)$$

where α_i are weights determined by a least square minimization principle and PR_i are the different candidate parameters. Time rate of change

THANKS
